

Counting statistics of single electron transport in a quantum dot



Simon Gustavsson

Renaud Leturcq

Barbara Simovič

R. Schleser

Thomas Ihn

Prof. Klaus Ensslin



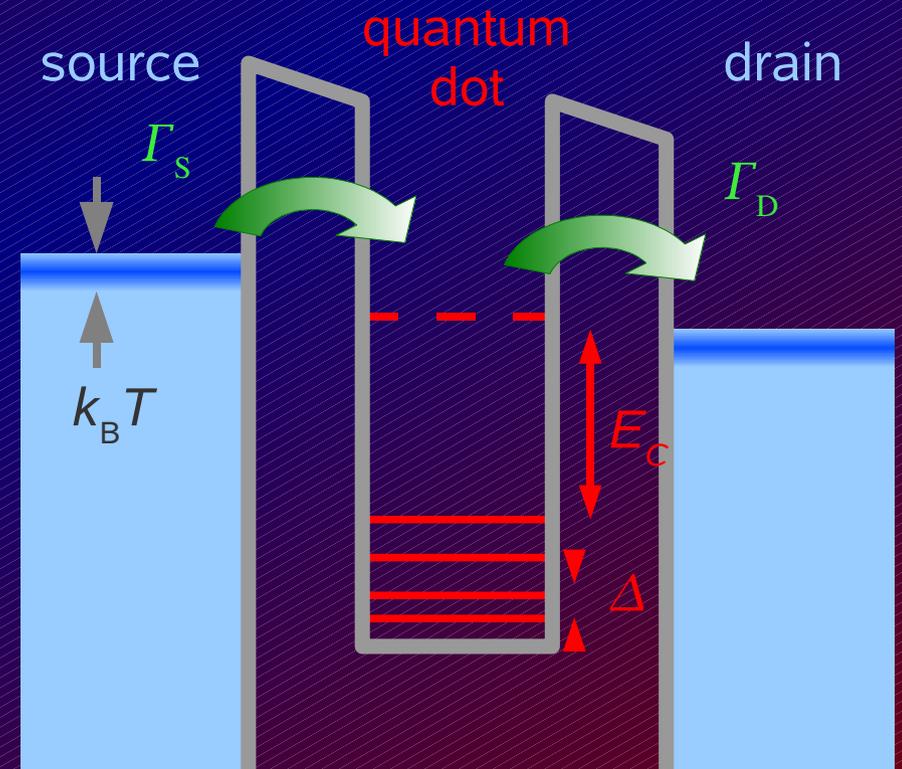
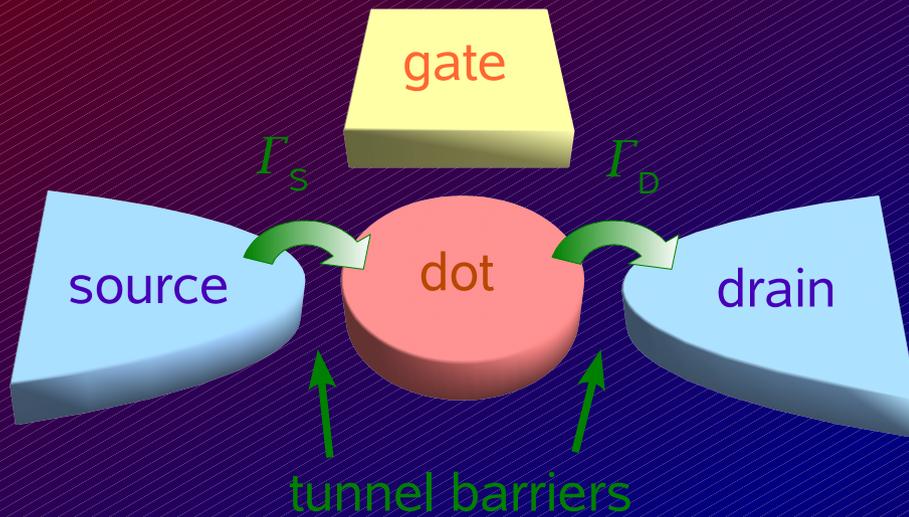
D. C. Driscoll
A. C. Gossard



Outline

- Introduction
- Detection of single electron transport
- Current fluctuations and full counting statistics in a semiconductor quantum dot
- Tunneling through multiple states: bunching
- Conclusion

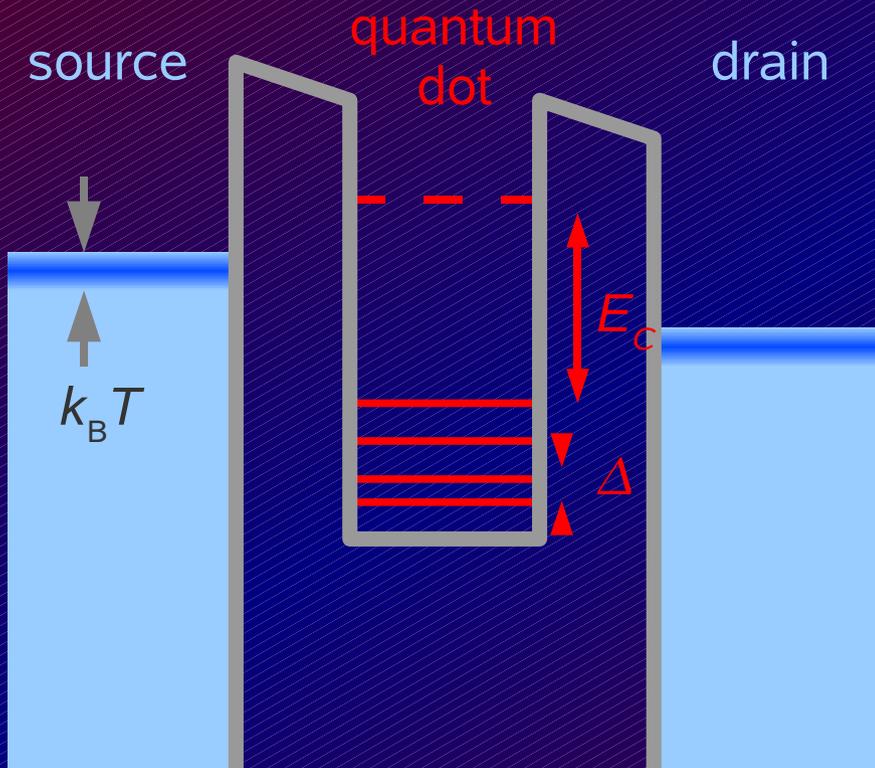
Semiconductor quantum dots



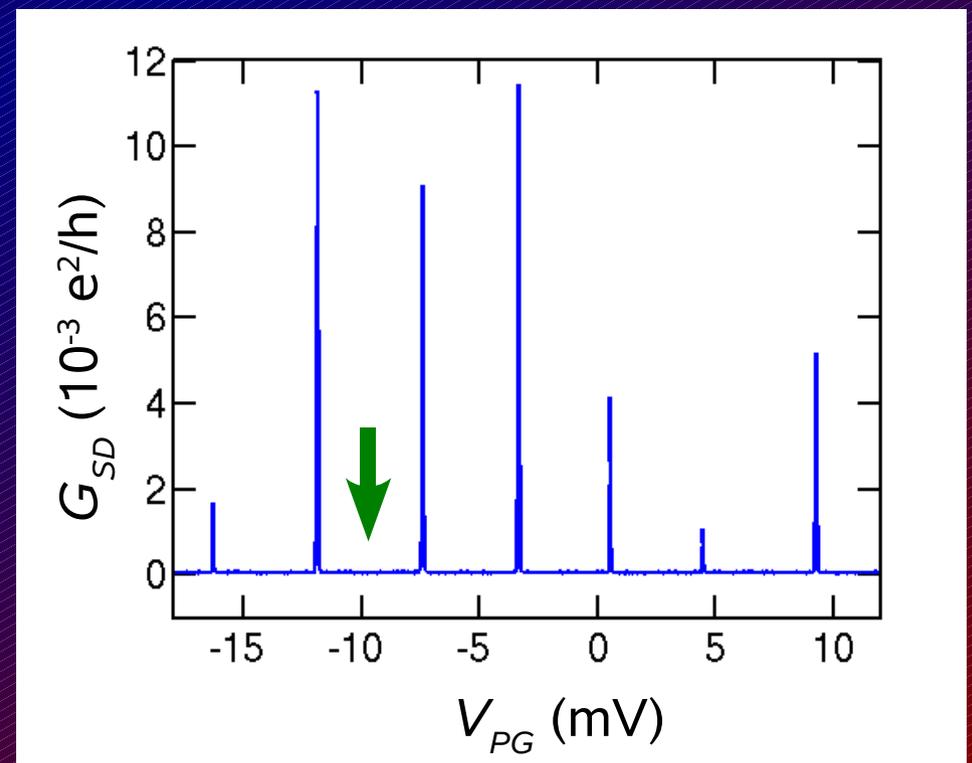
- Electrostatic energy E_C
- Quantum level spacing Δ

DC current measurement in a quantum dot

- Spectroscopy of electronic states

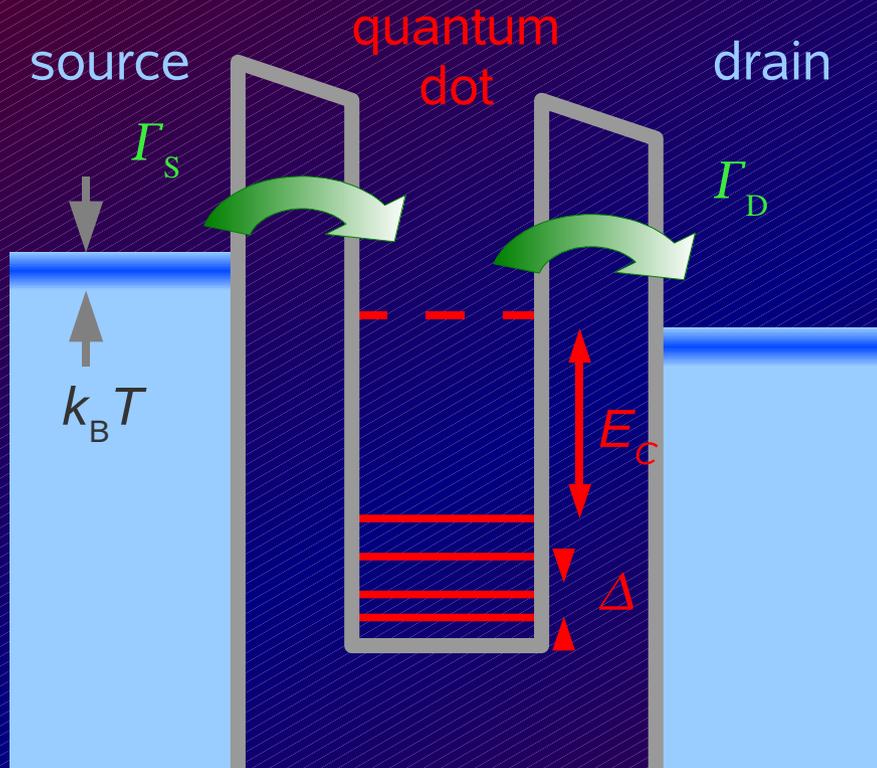


$$k_B T \ll \Delta < E_C$$

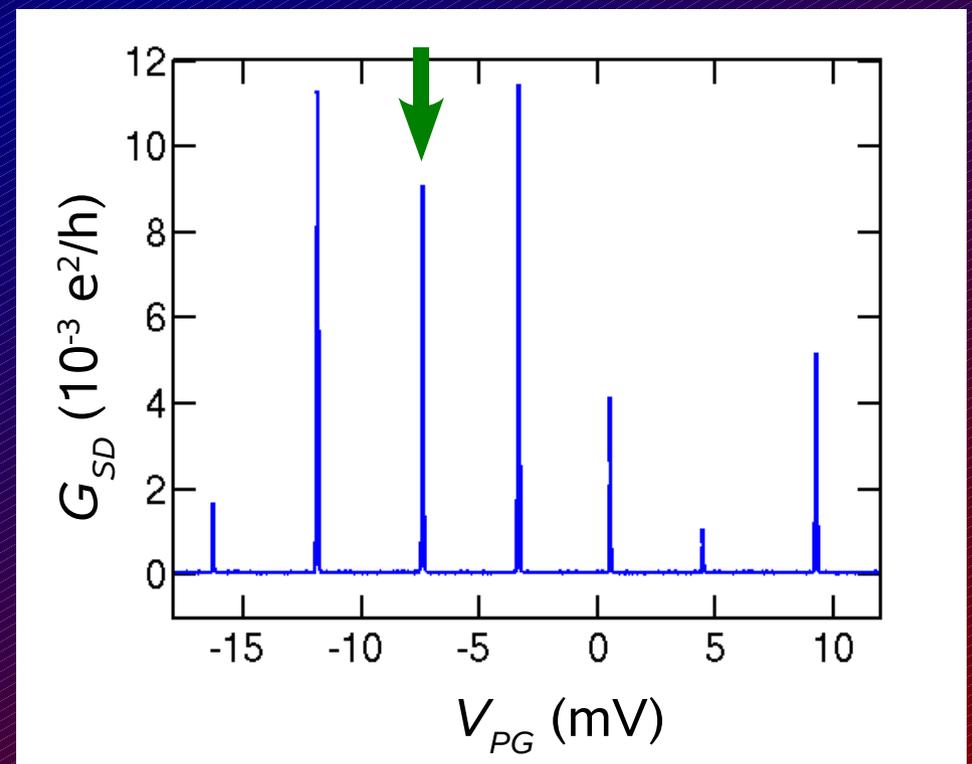


DC current measurement in a quantum dot

- Spectroscopy of electronic states

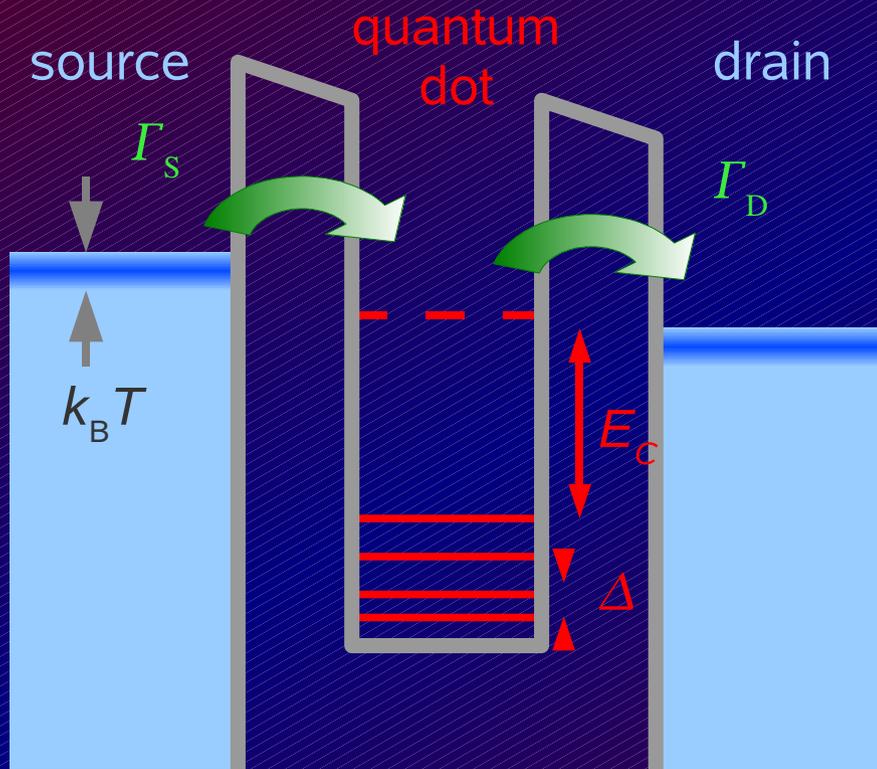


$$k_B T \ll \Delta < E_C$$

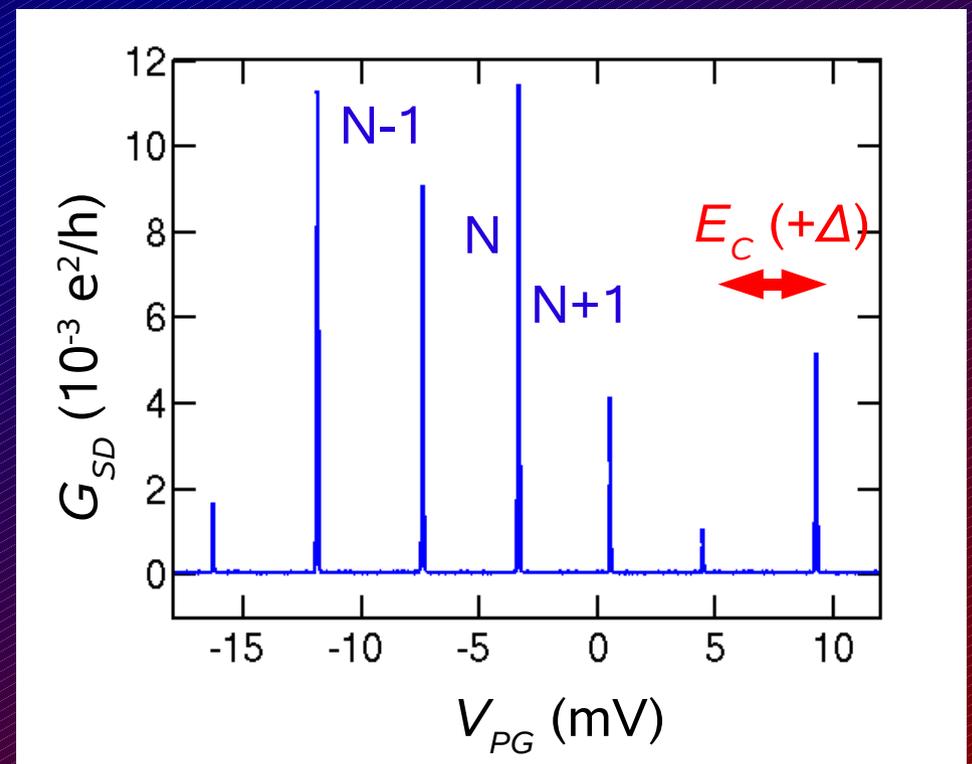


DC current measurement in a quantum dot

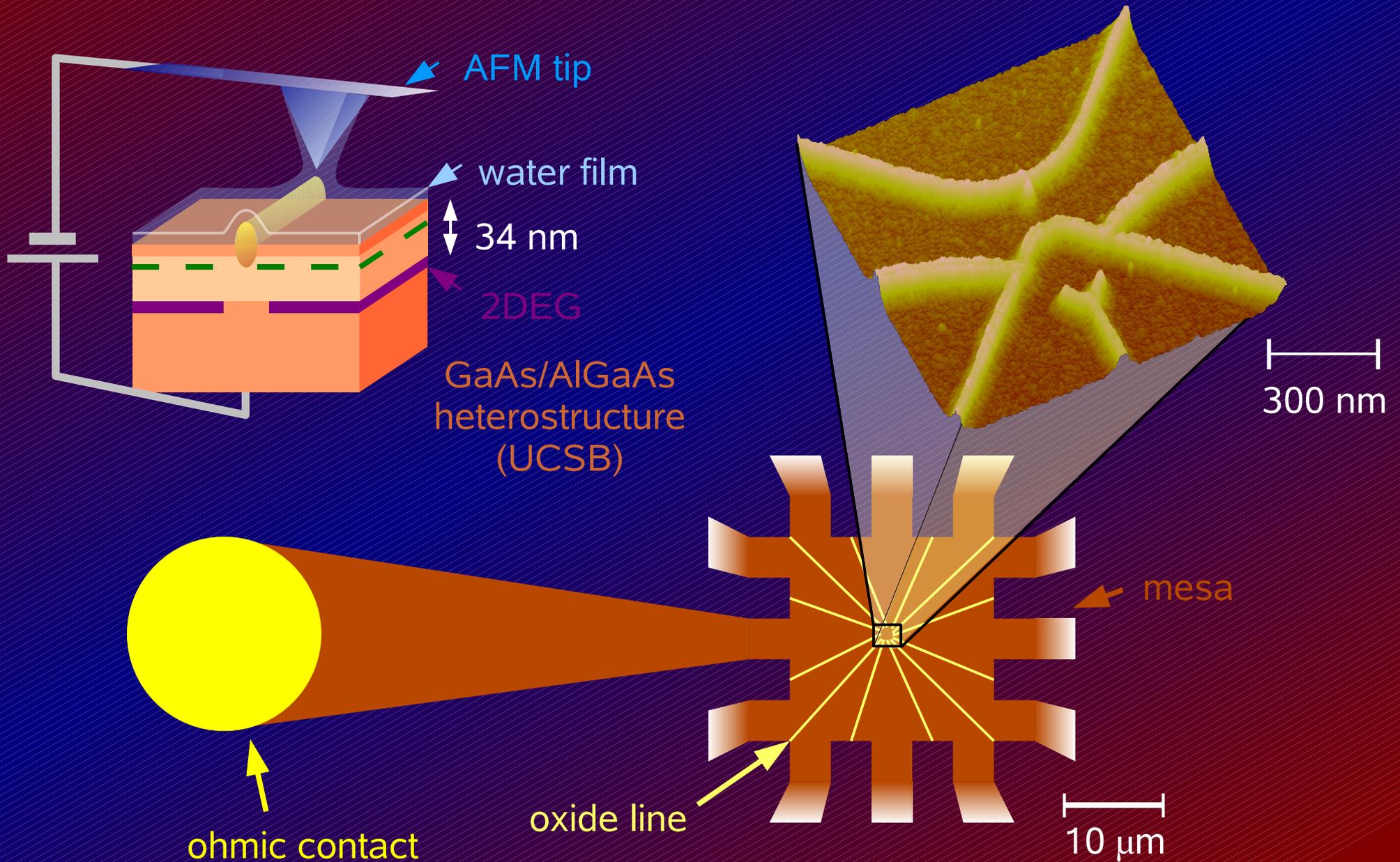
- Spectroscopy of electronic states



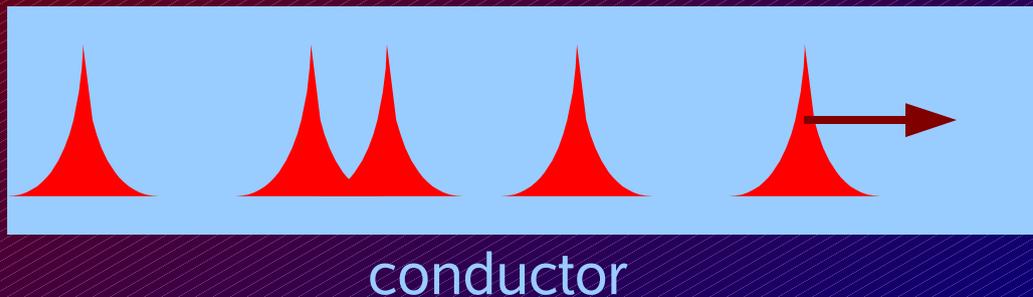
$$k_B T \ll \Delta < E_C$$



Quantum dots realized by AFM lithography



Measurement of current fluctuations



- Shot noise due to discreteness of charges
 - classical shot noise for independent particles (Poissonian noise): $S_I = 2eI$
- Usual measurement limited by noise of the current-meter $\Rightarrow S_I^{\min} \approx 10^{-29} \text{ A}^2/\text{Hz}$

Noise in quantum dots

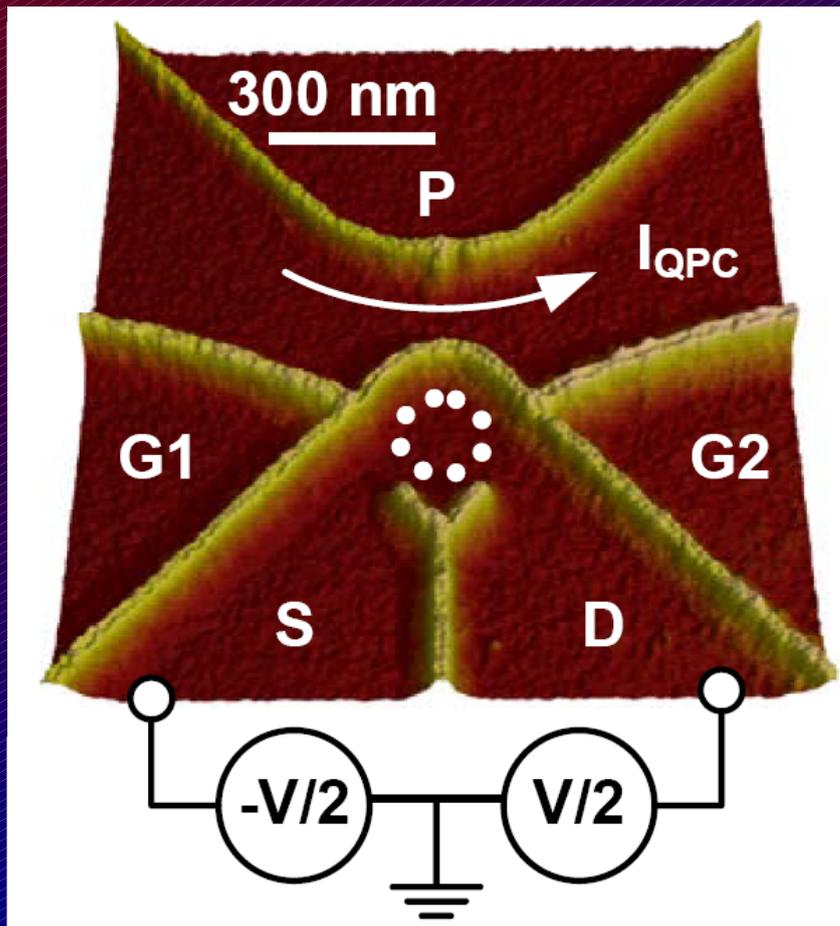
- Sub- vs. super-Poissonian shot noise in quantum dots
- Noise correlations in multi-terminal quantum dots
- Probing entanglement with the shot noise in quantum dot systems
- Kondo effect, spin blockade, etc...

Noise in quantum dots

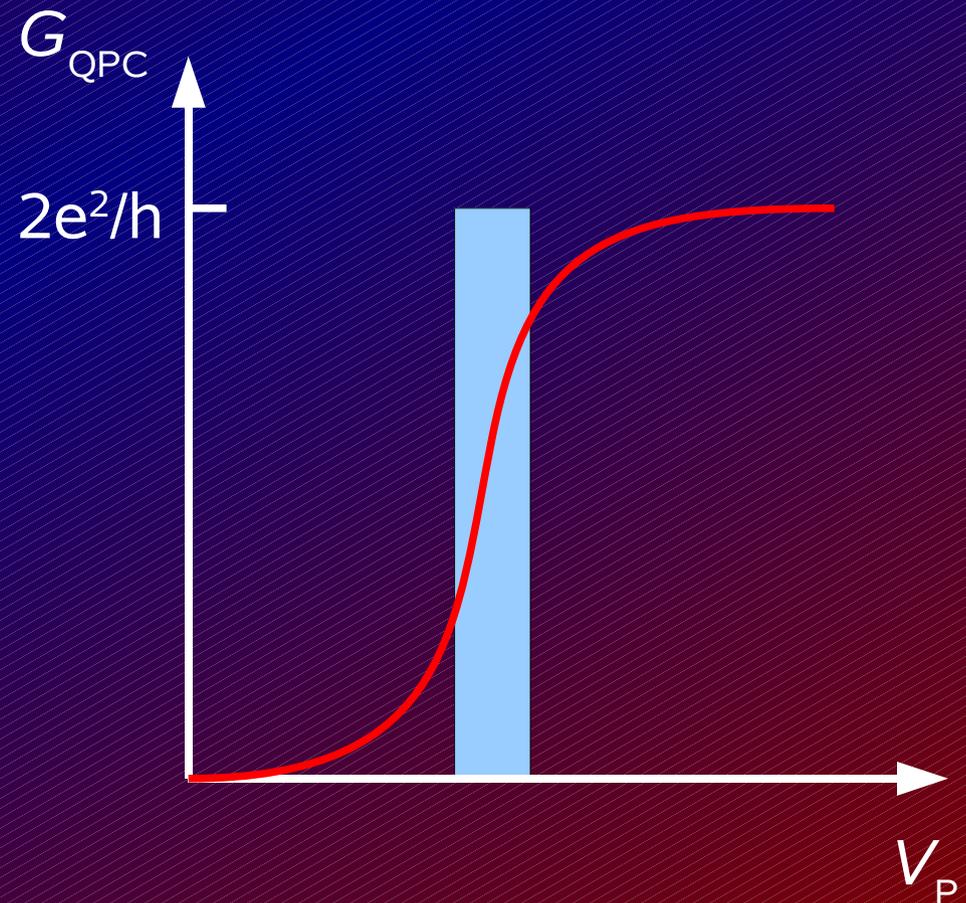
- Noise in interacting systems
 - deviations from Poissonian shot noise
- Early experiments in non-tunable quantum dots showed reduction of the shot noise: $S_I < 2eI$
 - Birk *et al.*, PRL **75**, 1610 (1995)
 - Nauen *et al.*, PRB **70**, 033305 (2004)
- Challenge in lateral quantum dots
 - very low noise level: $I < 1$ pA $\Rightarrow S_I < 10^{-31}$ A²/Hz !
 - strongly non-linear systems

Detection of single electron transport

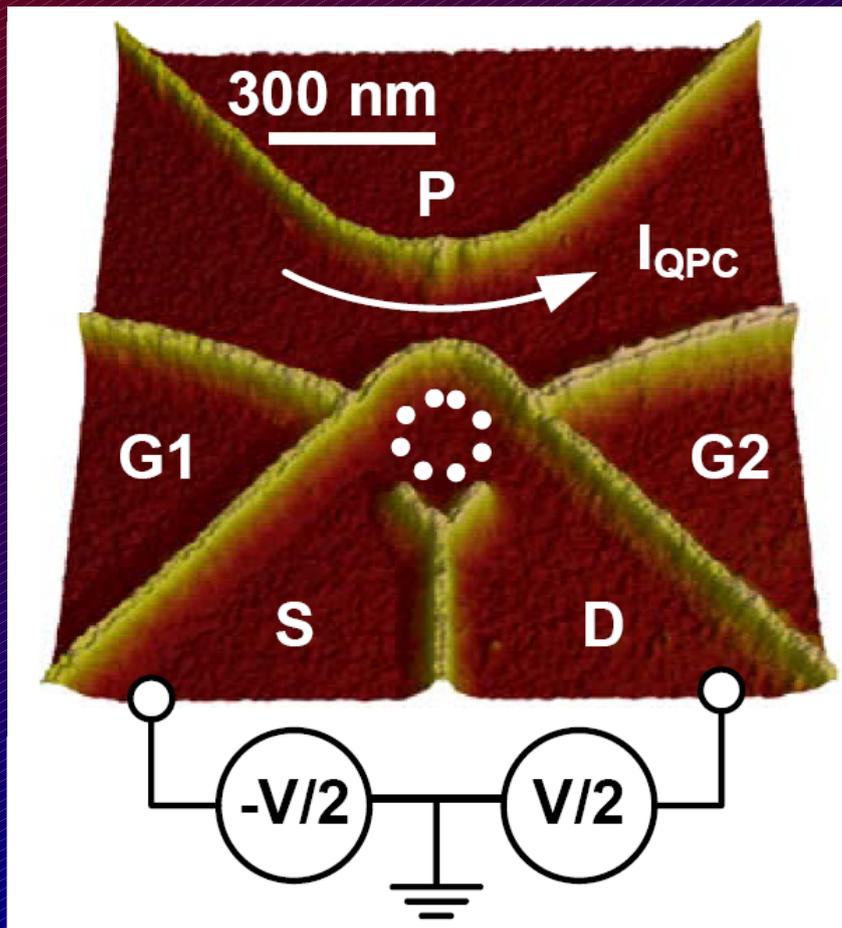
- Quantum point contact as a charge detector



$$T_e = 350 \text{ mK}$$

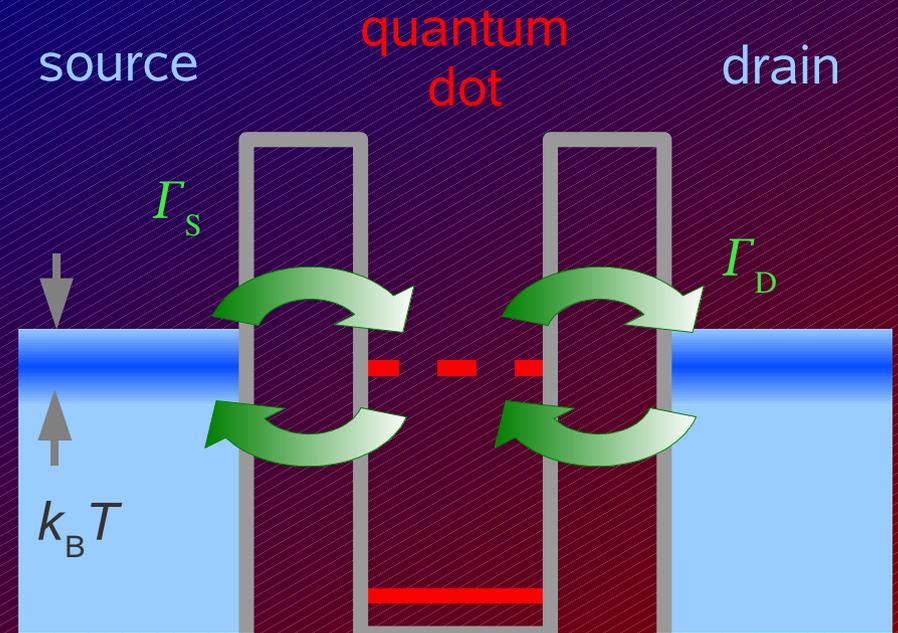


Detection of single electron transport

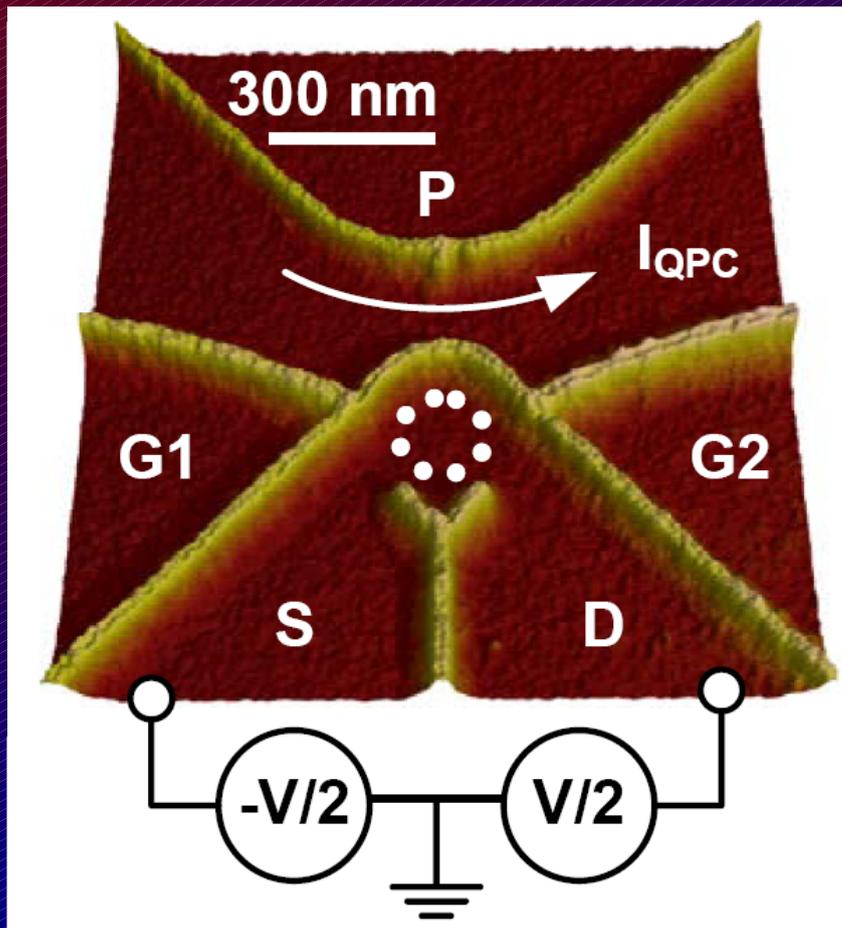


$$T_e = 350 \text{ mK}$$

- Quantum point contact as a charge detector
- Low bias voltage on the quantum dot

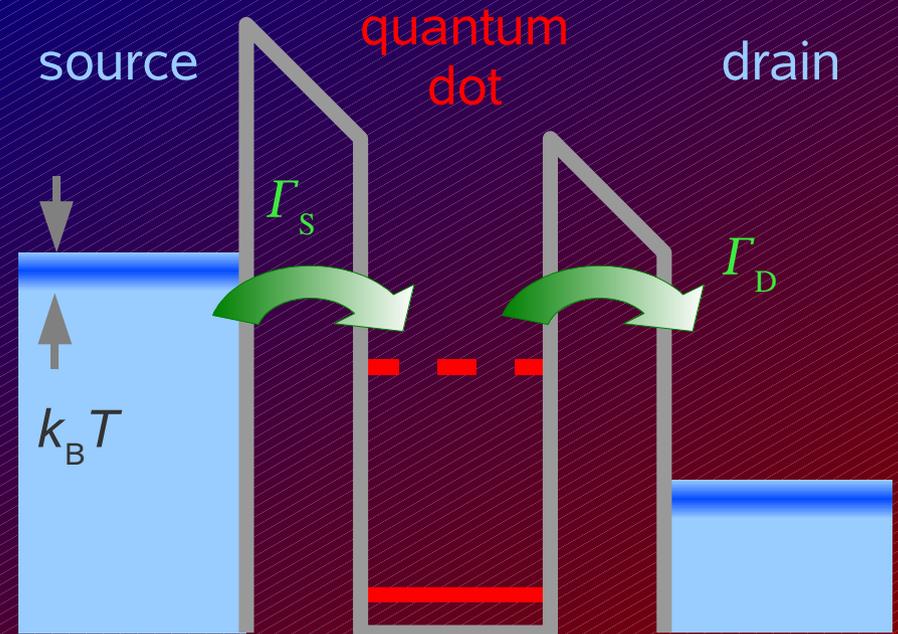


Detection of single electron transport

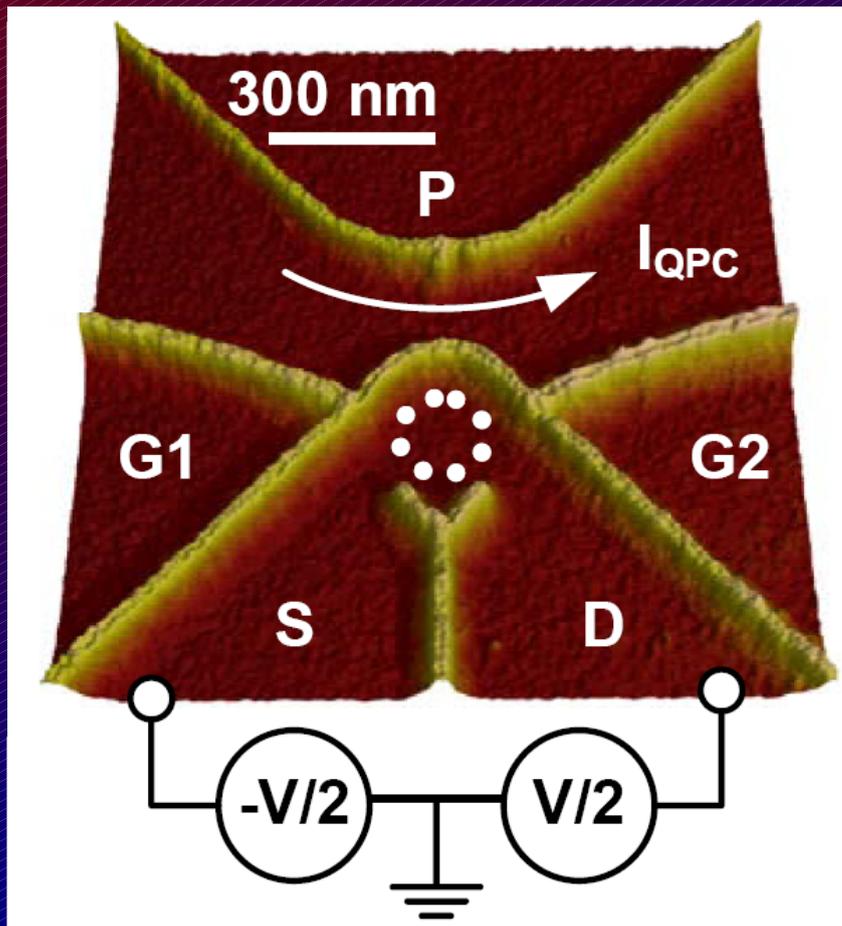


$$T_e = 350 \text{ mK}$$

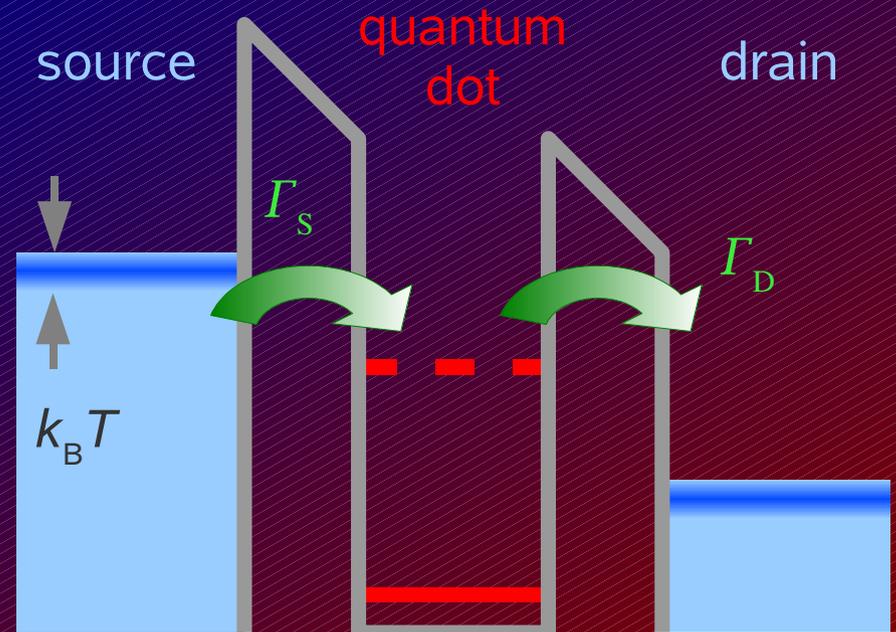
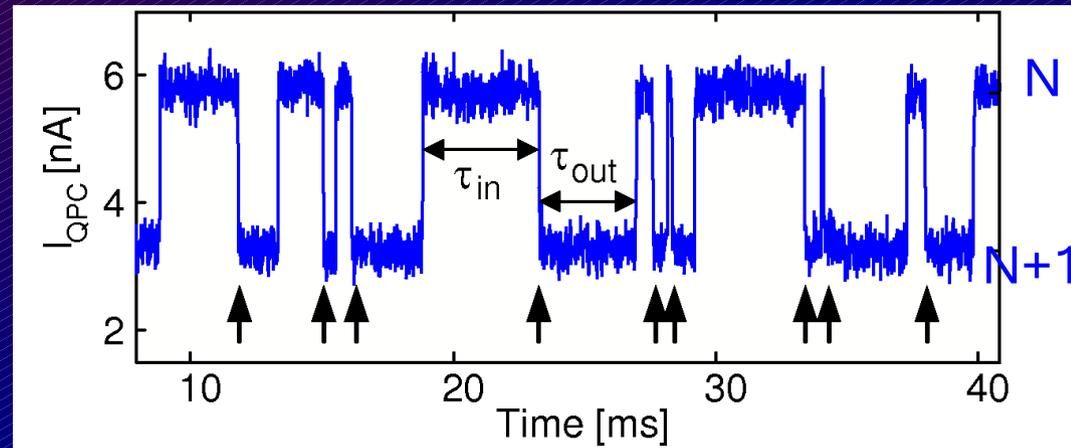
- Quantum point contact as a charge detector
- Large bias voltage on the quantum dot



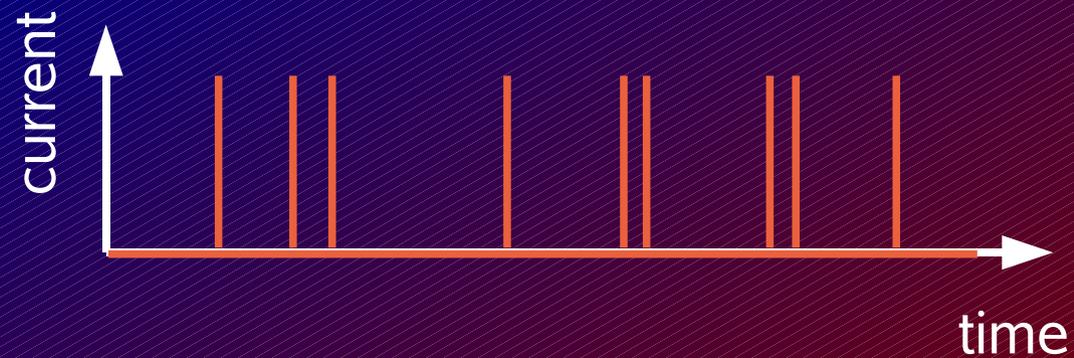
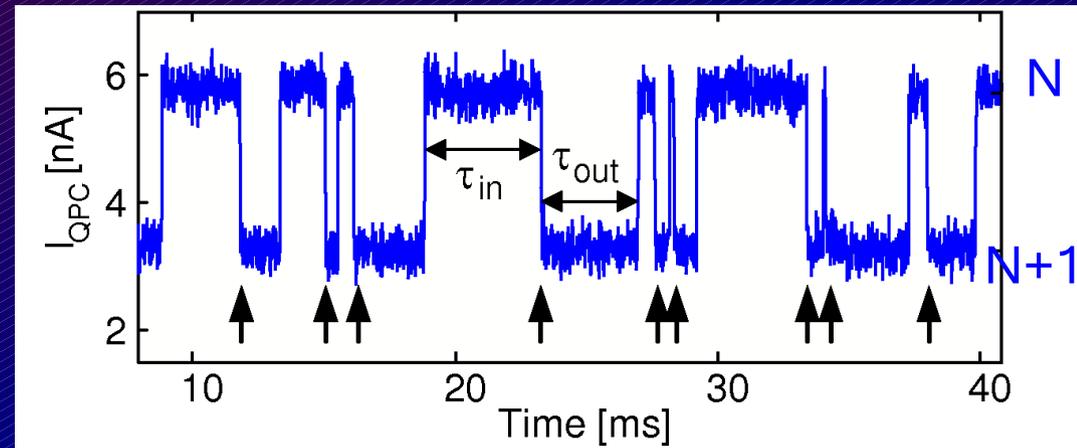
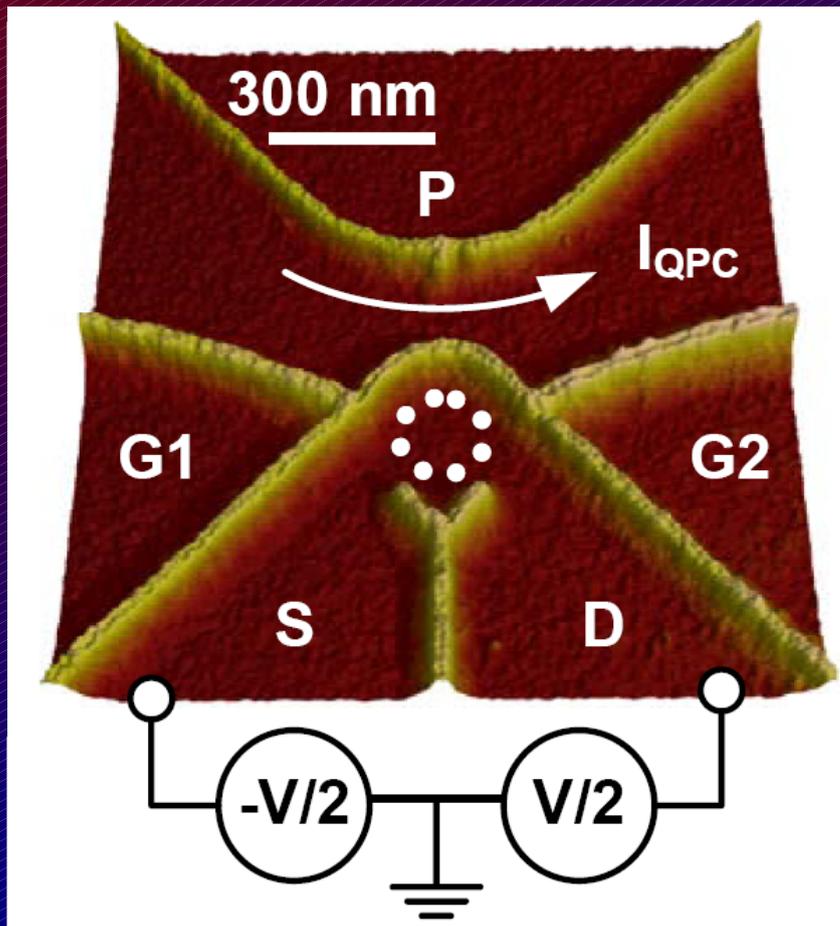
Detection of single electron transport



$$T_e = 350 \text{ mK}$$



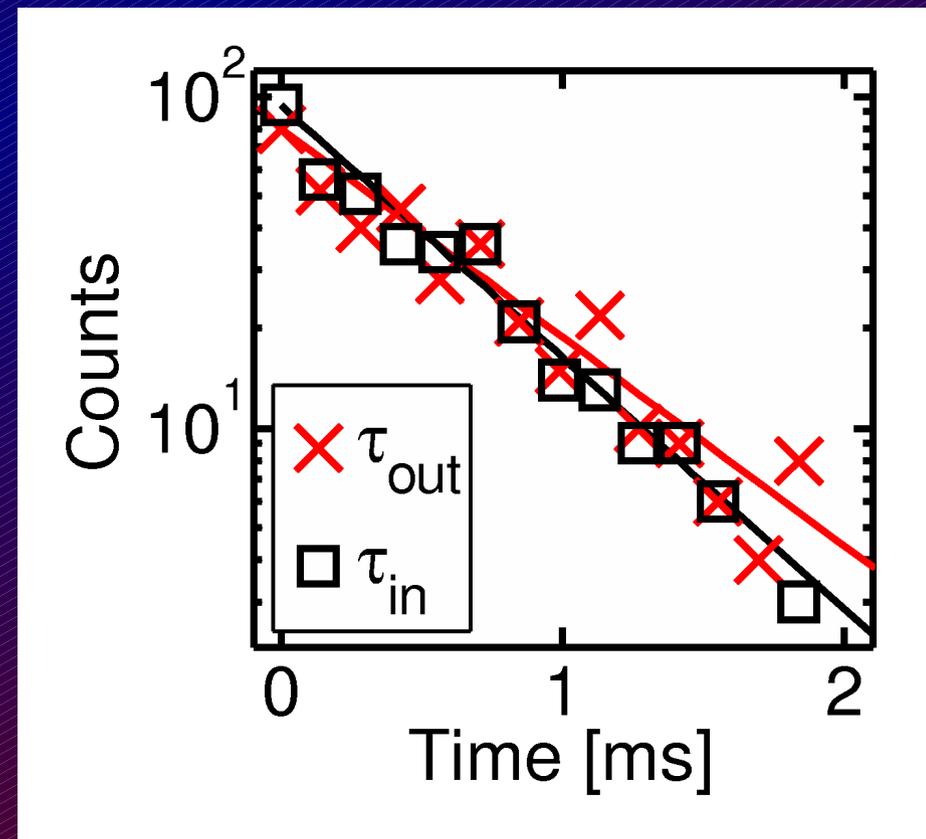
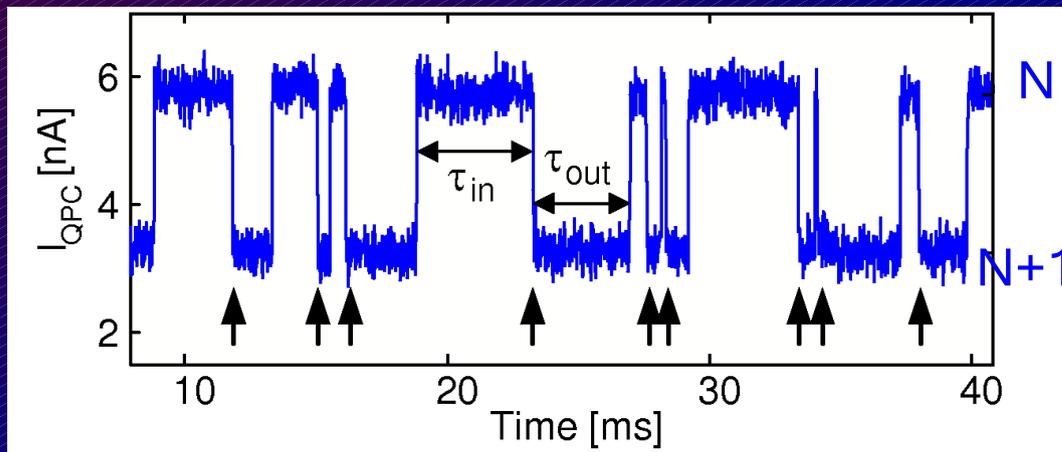
Detection of single electron transport



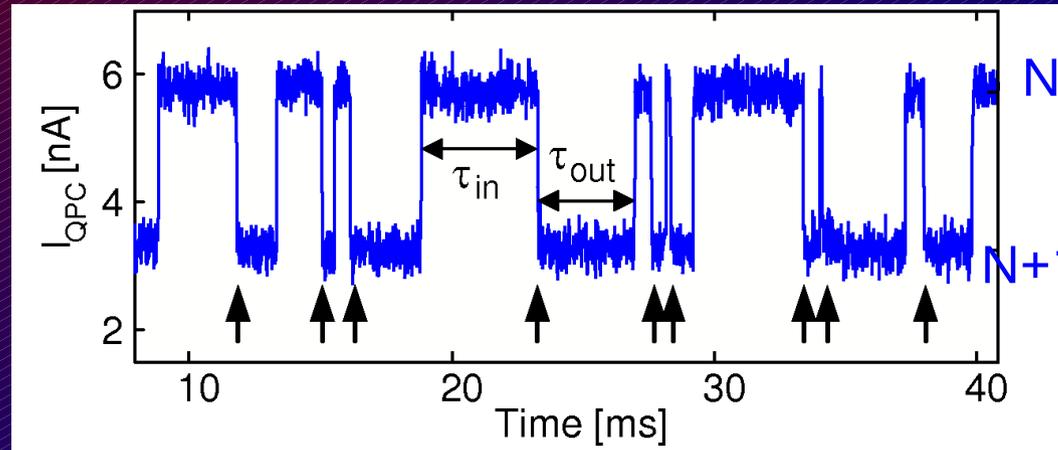
$$T_e = 350 \text{ mK}$$

Determination of the individual tunneling rates

- Exponential distribution of waiting times for independent events
- $\Gamma_S = \langle \tau_{in} \rangle$, $\Gamma_D = \langle \tau_{out} \rangle$



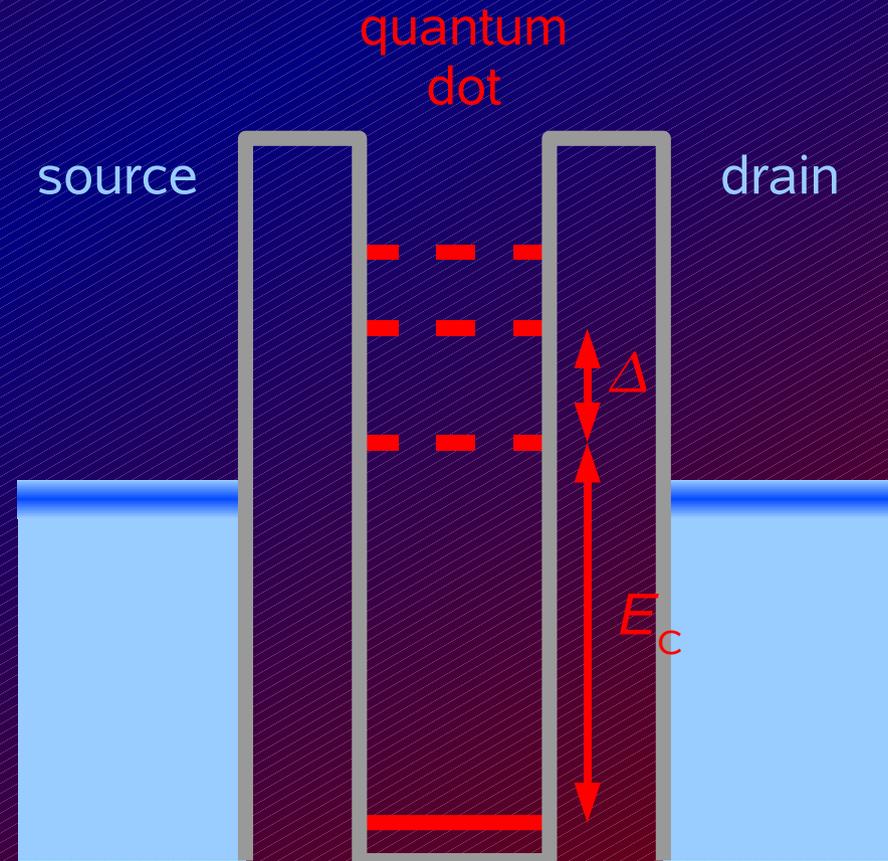
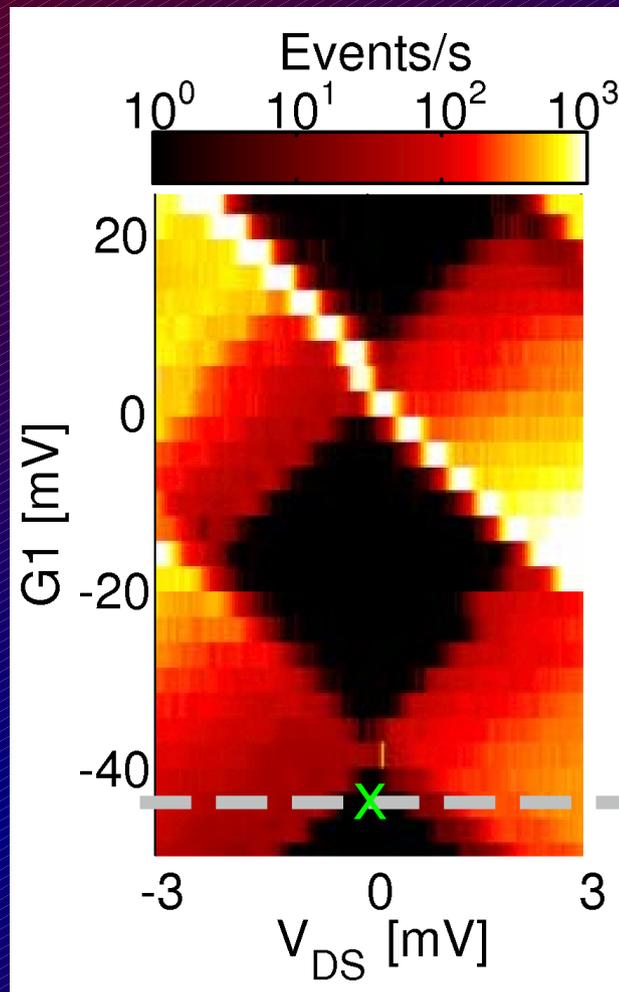
Measuring the current by counting electrons



- Count number n of electron entering the dot within a time t_0 : $I = e\langle n \rangle / t_0$
- Max. current = few fA (bandwidth = 30 kHz)
- BUT no absolute limitation for low current and noise measurements
 - we show here: $I \approx \text{few aA}$, $S_I \approx 10^{-35} \text{ A}^2/\text{Hz}$

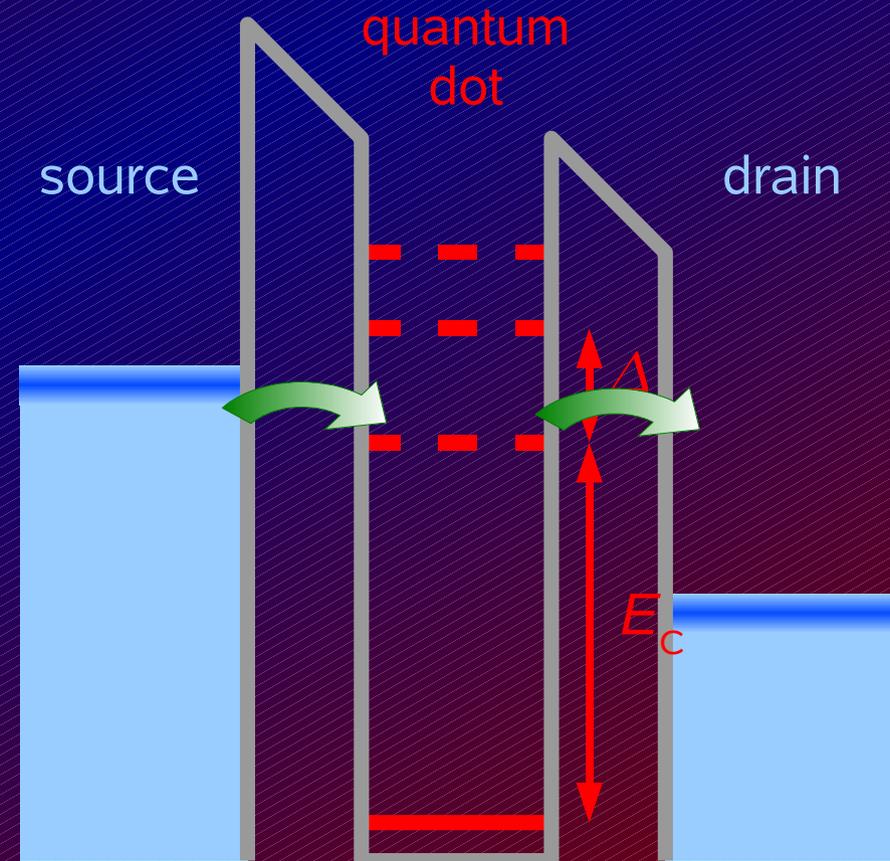
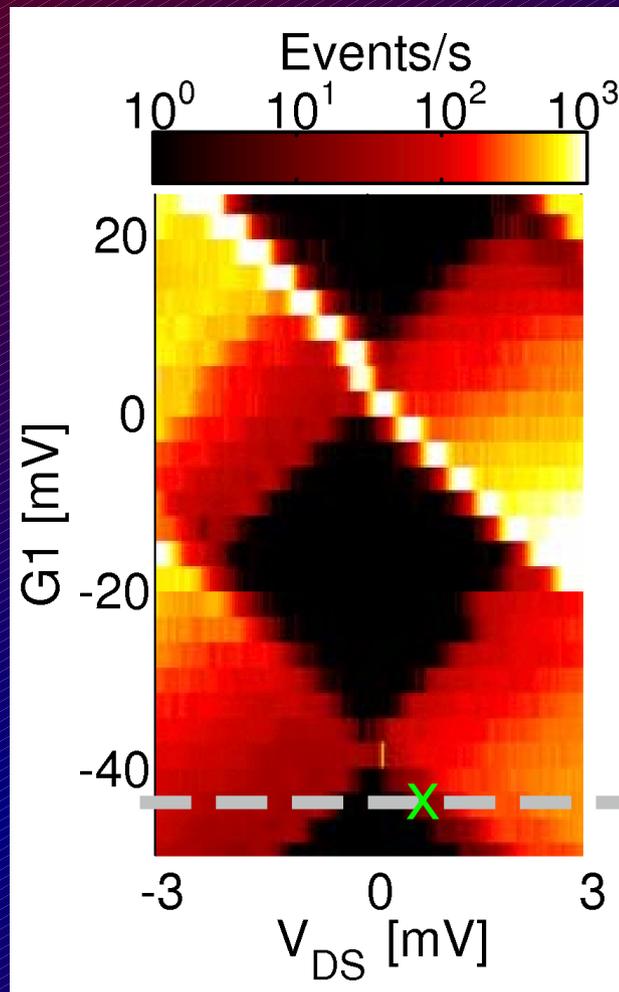
Coulomb diamond measured by electron counting

➤ $I = e\langle n \rangle / t_0$



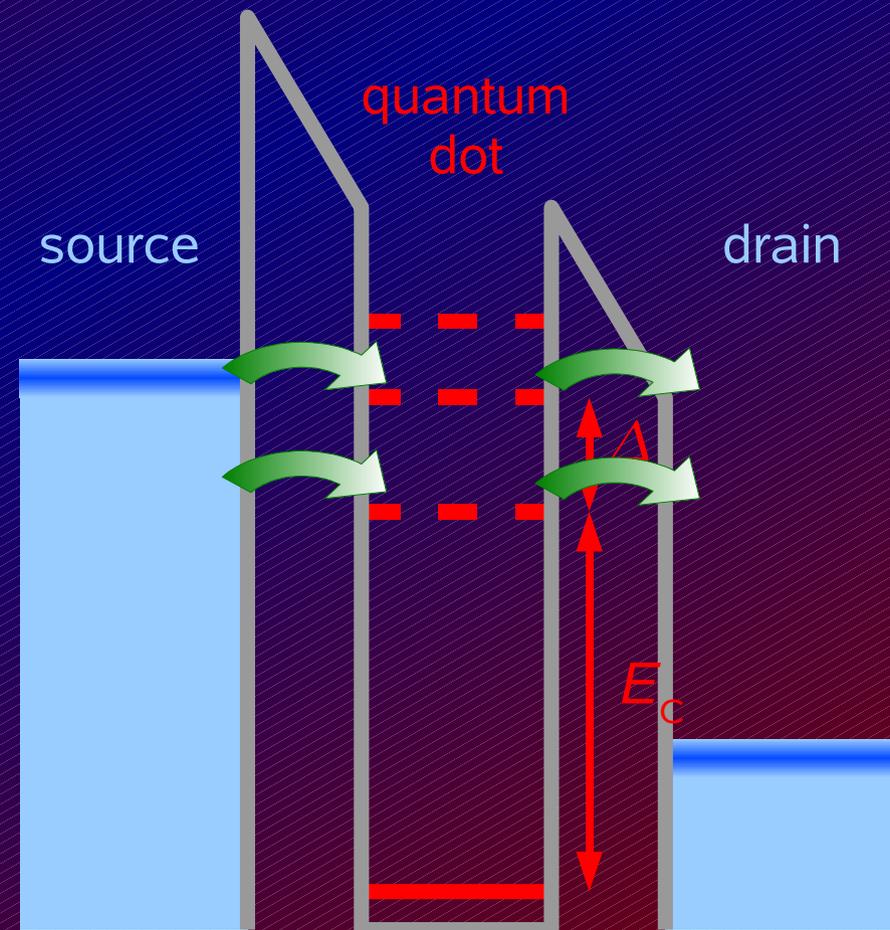
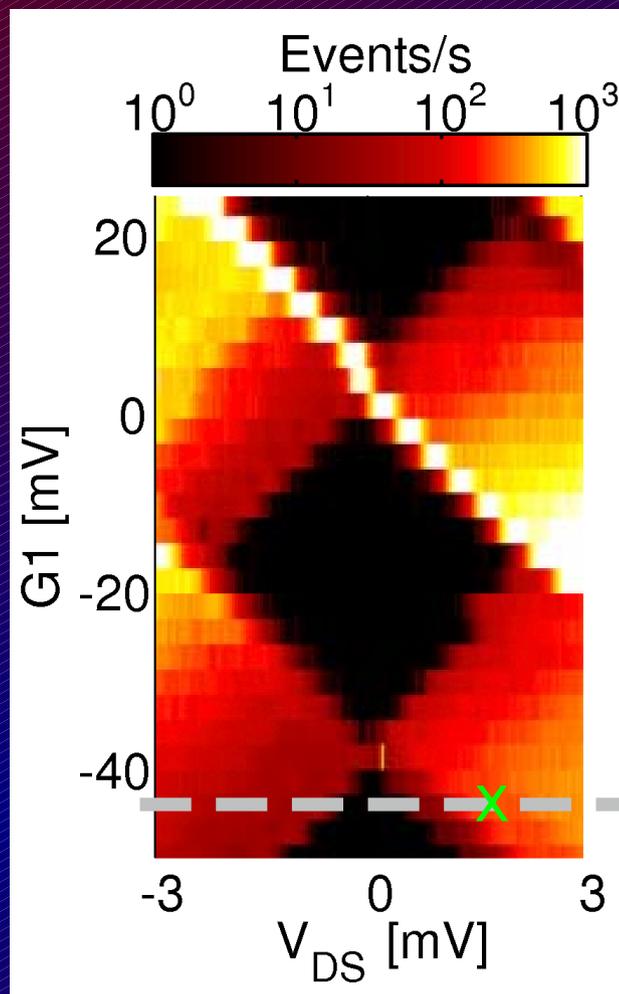
Coulomb diamond measured by electron counting

➤ $I = e\langle n \rangle / t_0$



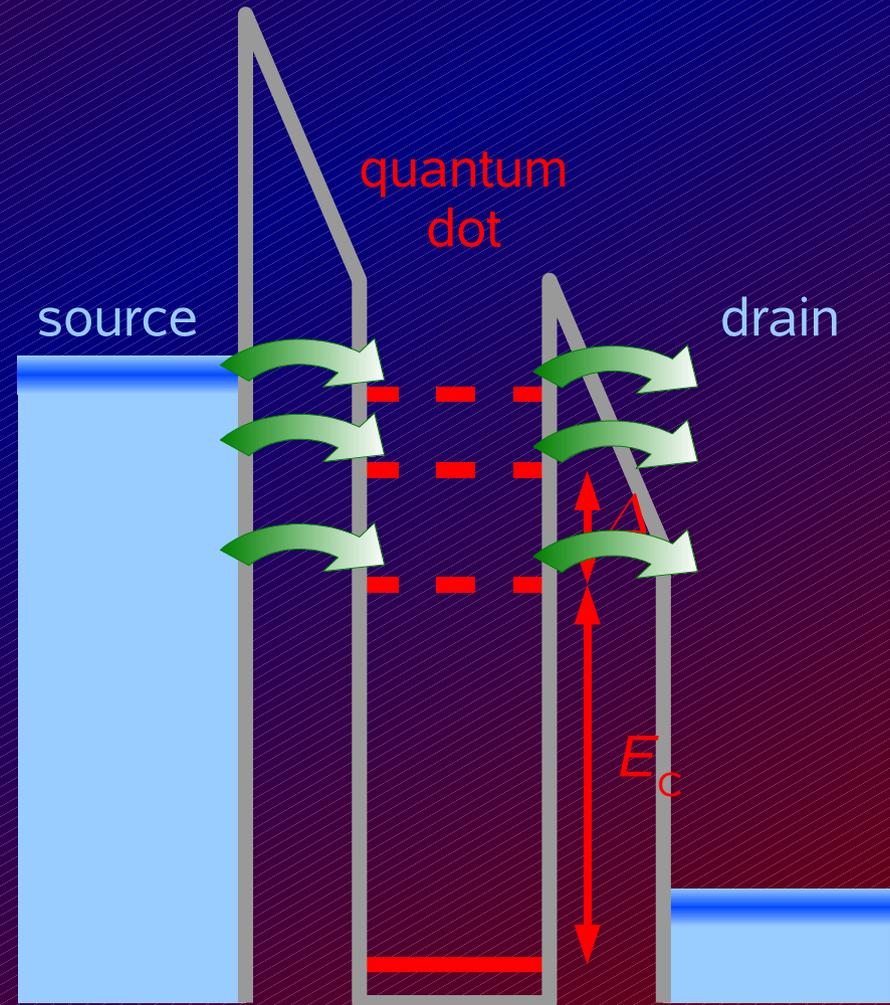
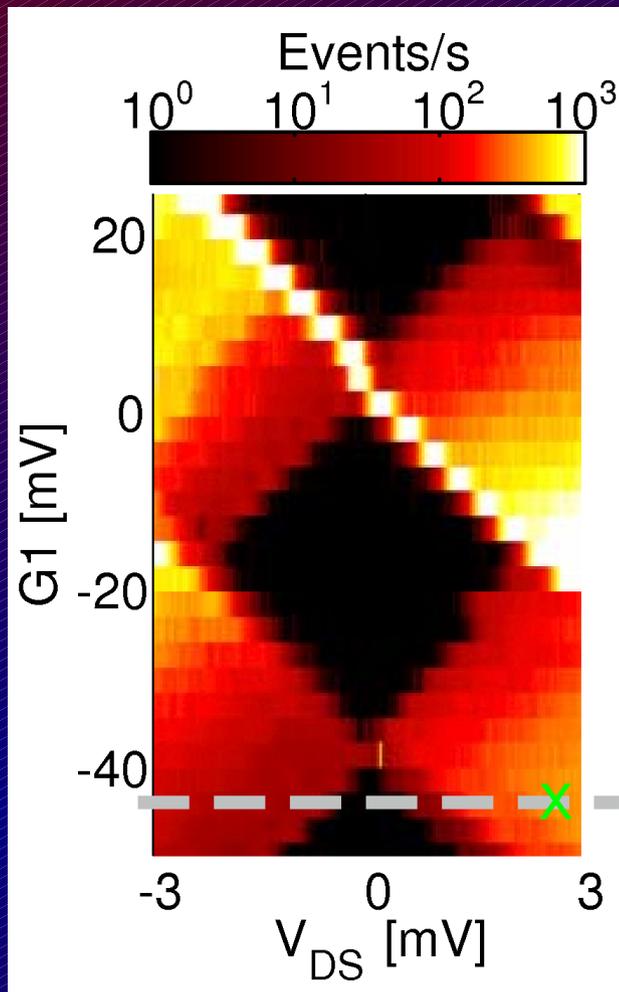
Coulomb diamond measured by electron counting

➤ $I = e\langle n \rangle/t_0$



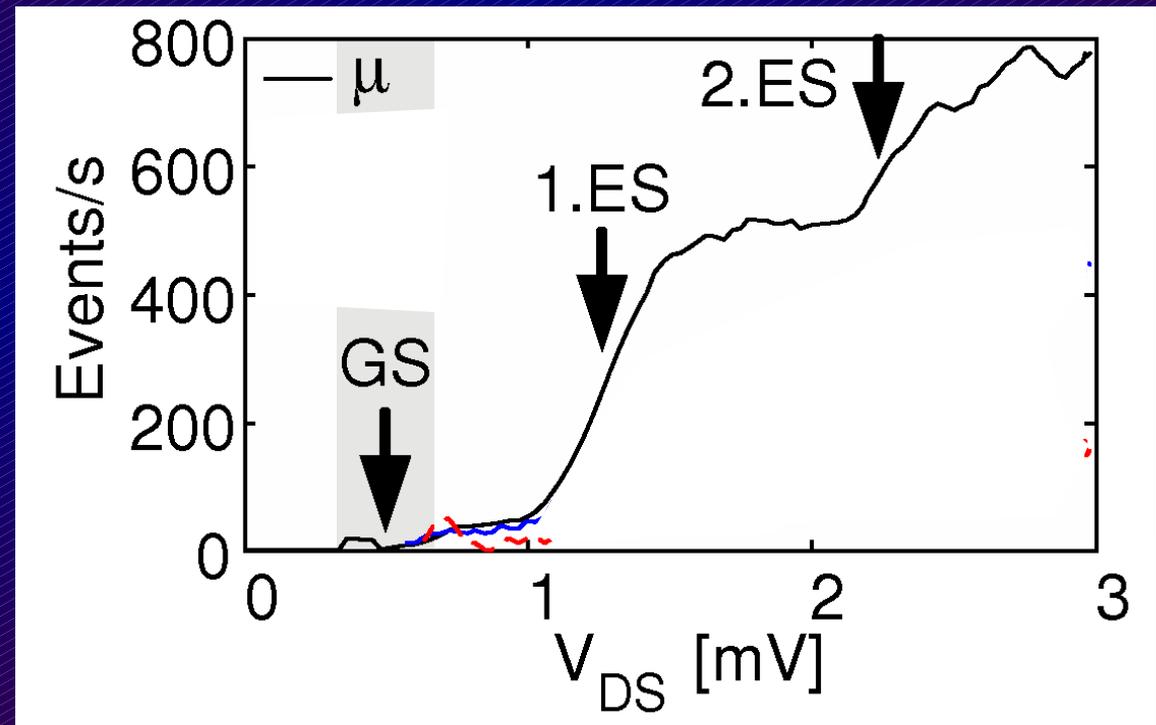
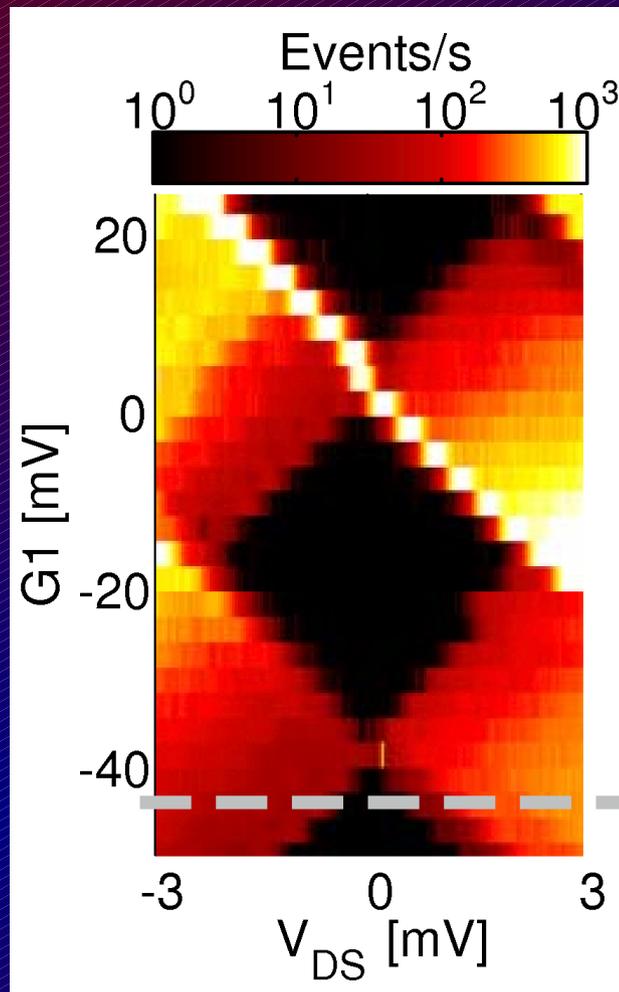
Coulomb diamond measured by electron counting

➤ $I = e\langle n \rangle/t_0$



Coulomb diamond measured by electron counting

➤ $I = e\langle n \rangle / t_0$



Current fluctuations measured by electron counting

- More than noise: access to the full counting statistics (distribution function)

- $I = e\mu/t_0,$

- $\mu = \langle n \rangle$

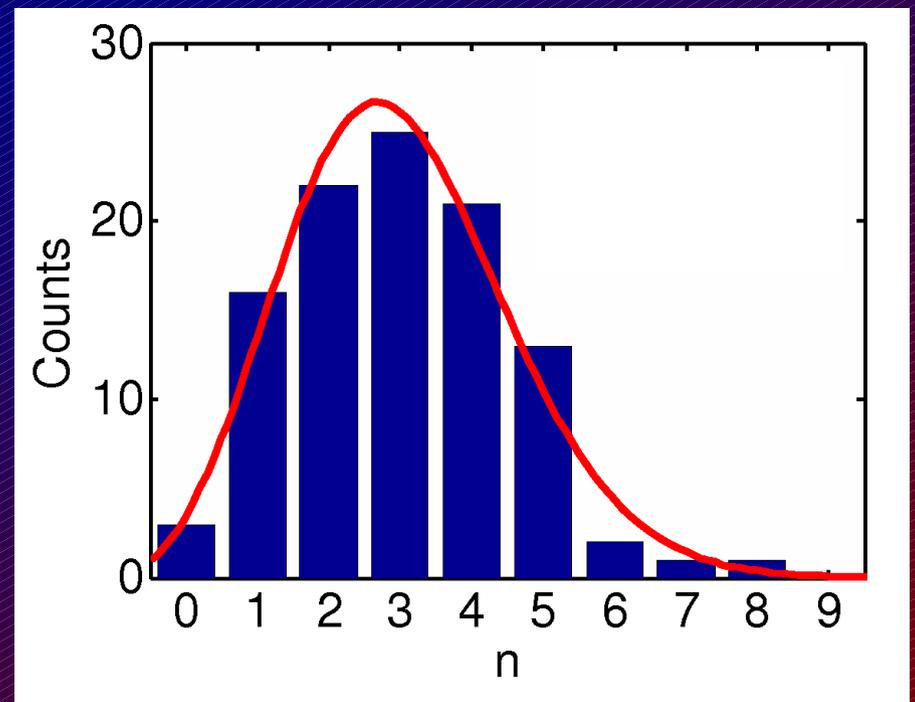
- $S_I = 2e^2\mu_2/t_0,$

- $\mu_2 = \langle (n - \langle n \rangle)^2 \rangle$

- $S_I^3 = e^3\mu_3/t_0,$

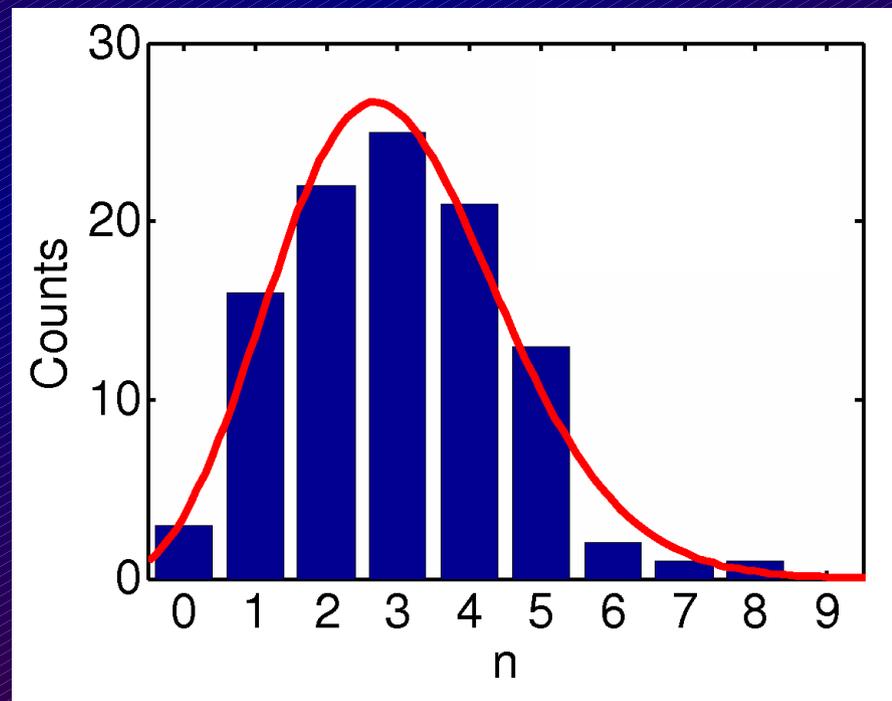
- $\mu_3 = \langle (n - \langle n \rangle)^3 \rangle$

- and many more...

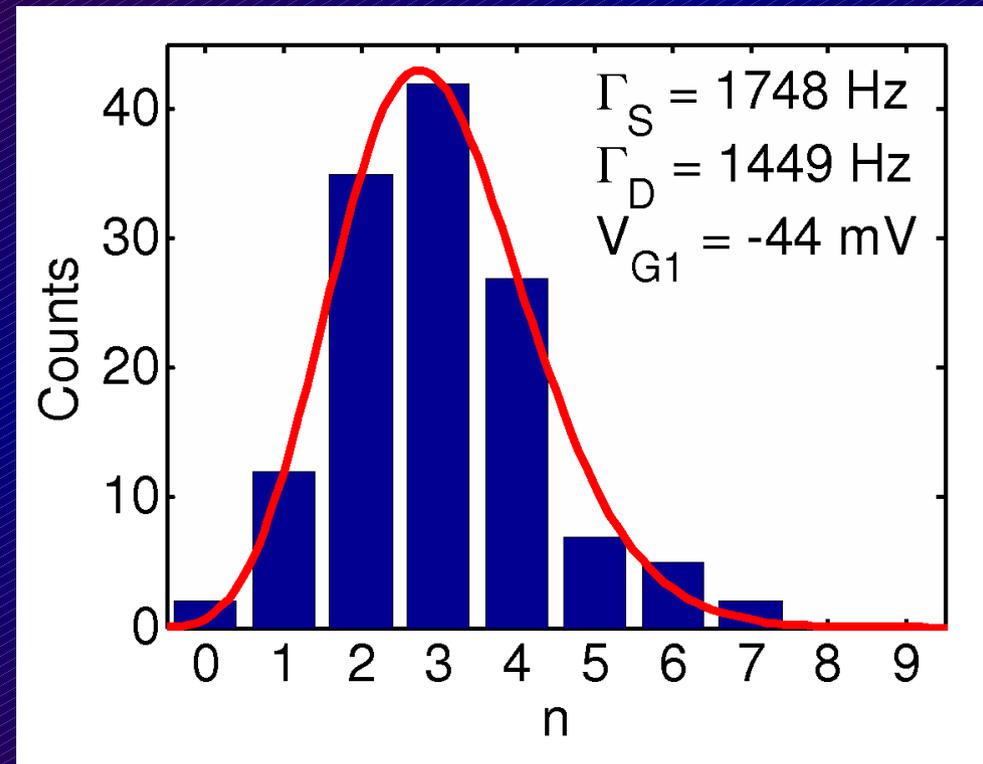
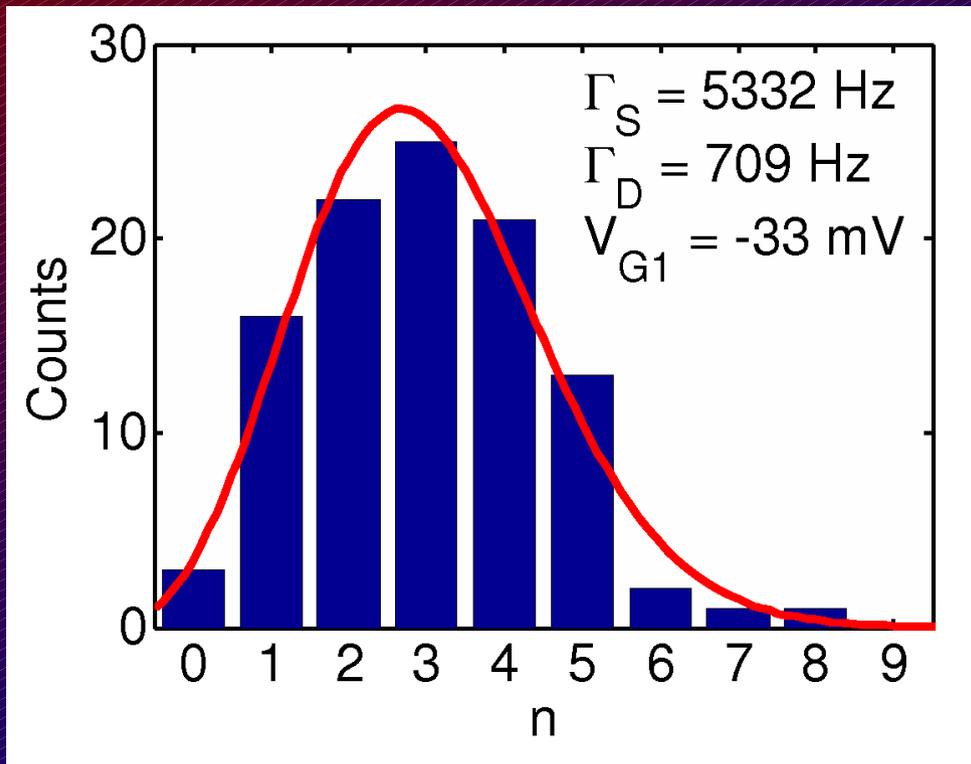


Distribution function for electrons in a conductor

- Classical noise for independent particles
⇒ Poisson distribution: $\mu = \mu_2 = \mu_3$
- Particles with repulsive interaction ⇒ sub-Poissonian distribution: $\mu_2 < \mu$, $\mu_3 < \mu, \dots$



Histogram of current fluctuations



➤ Poisson distribution for asymmetric coupling

➤ Sub-Poisson distribution for symmetric coupling

Theory: Hershfield *et al.*, PRB **47**, 1967 (1993)

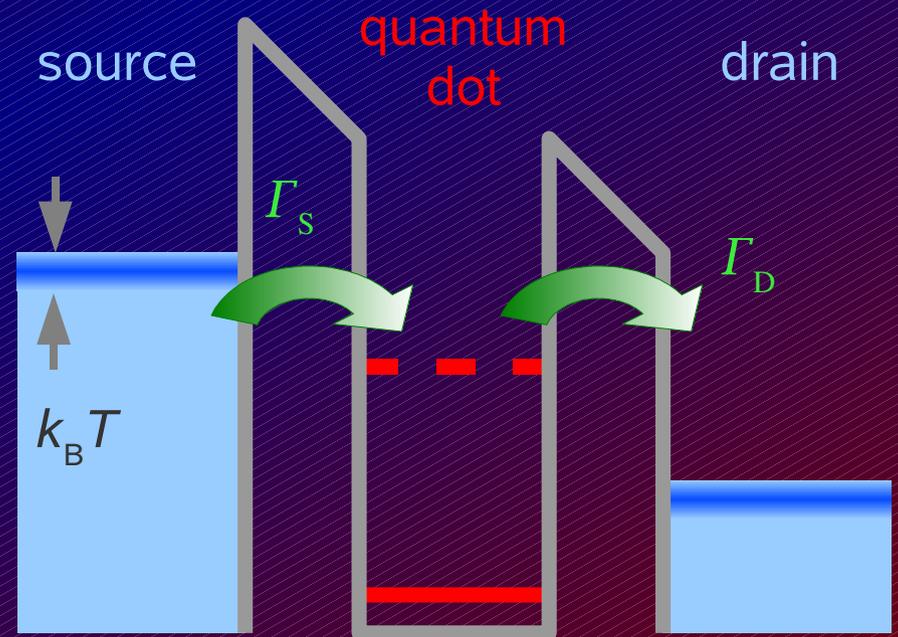
Bagrets & Nazarov, PRB **67**, 085316 (2003)

Counting statistics in a single-level quantum dot

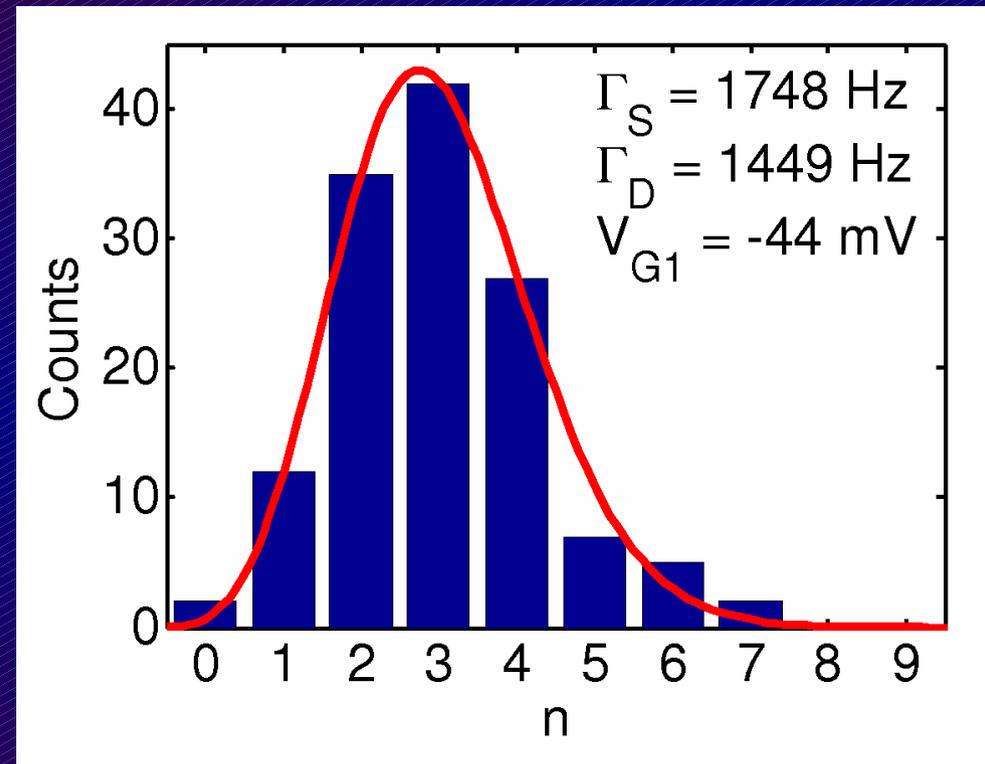
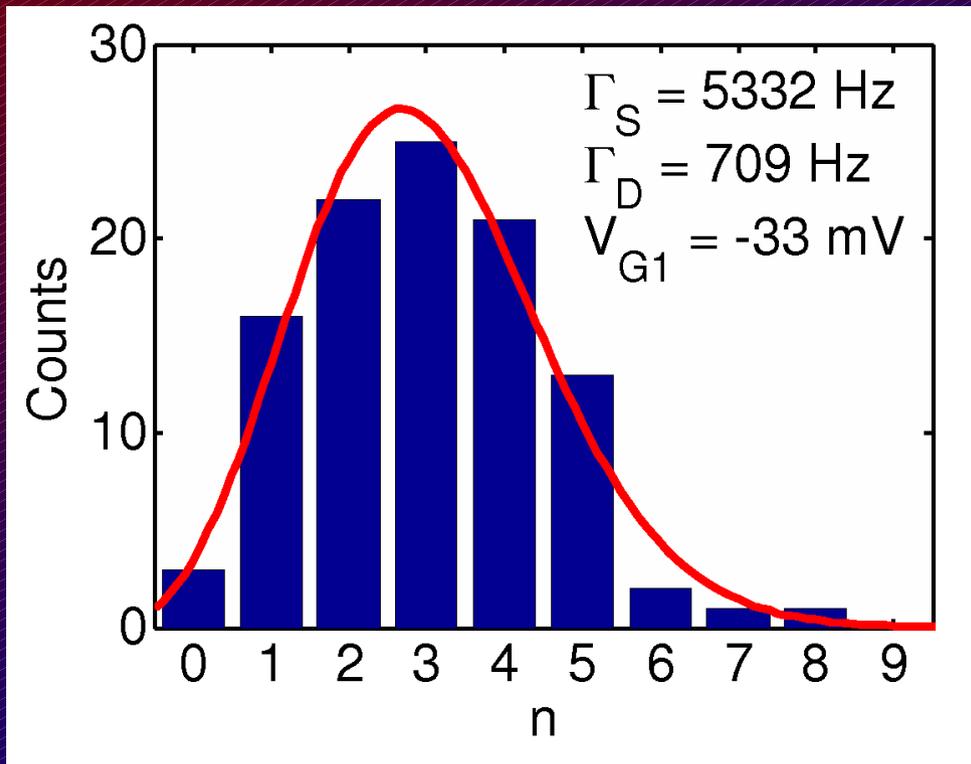
Bagrets & Nazarov, PRB **67**, 085316 (2003)

$$\frac{d}{dt} \begin{pmatrix} p_0 \\ p_1 \end{pmatrix} = M \begin{pmatrix} p_0 \\ p_1 \end{pmatrix}$$

$$M(\chi) = \begin{pmatrix} -\Gamma_D & \Gamma_D \\ \Gamma_S e^{i\chi} & -\Gamma_S \end{pmatrix}$$



Histogram of current fluctuations



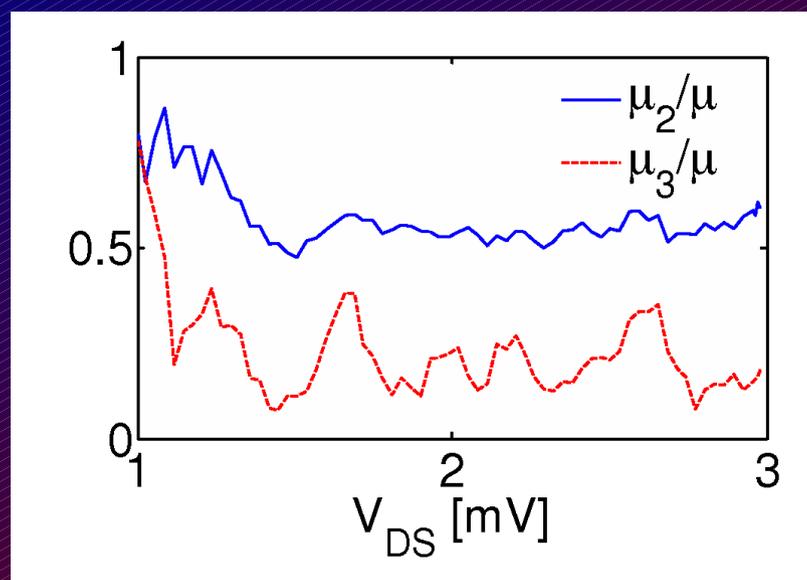
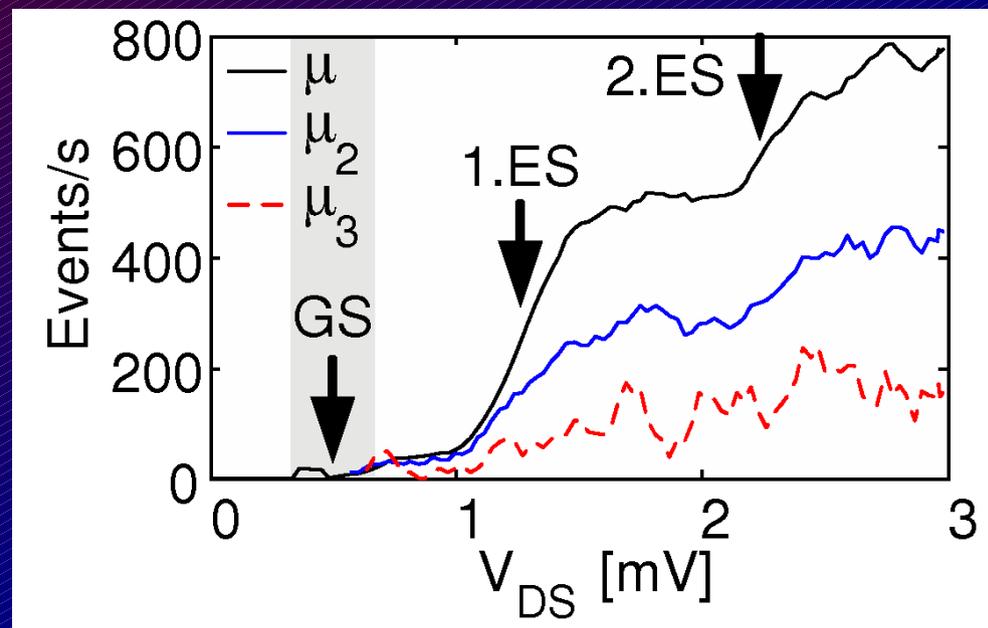
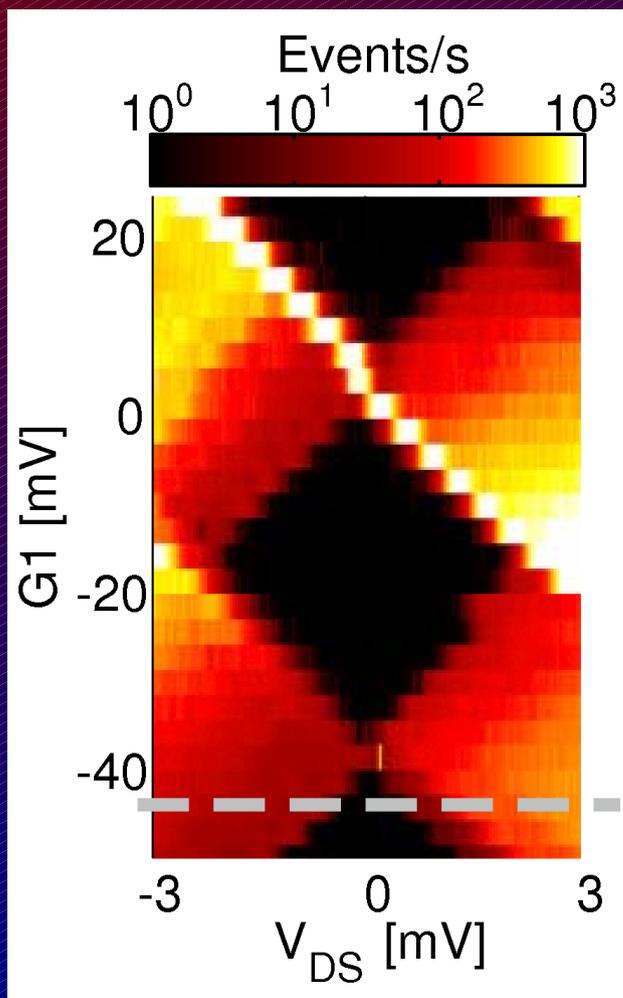
➤ Poisson distribution for asymmetric coupling

➤ Sub-Poisson distribution for symmetric coupling

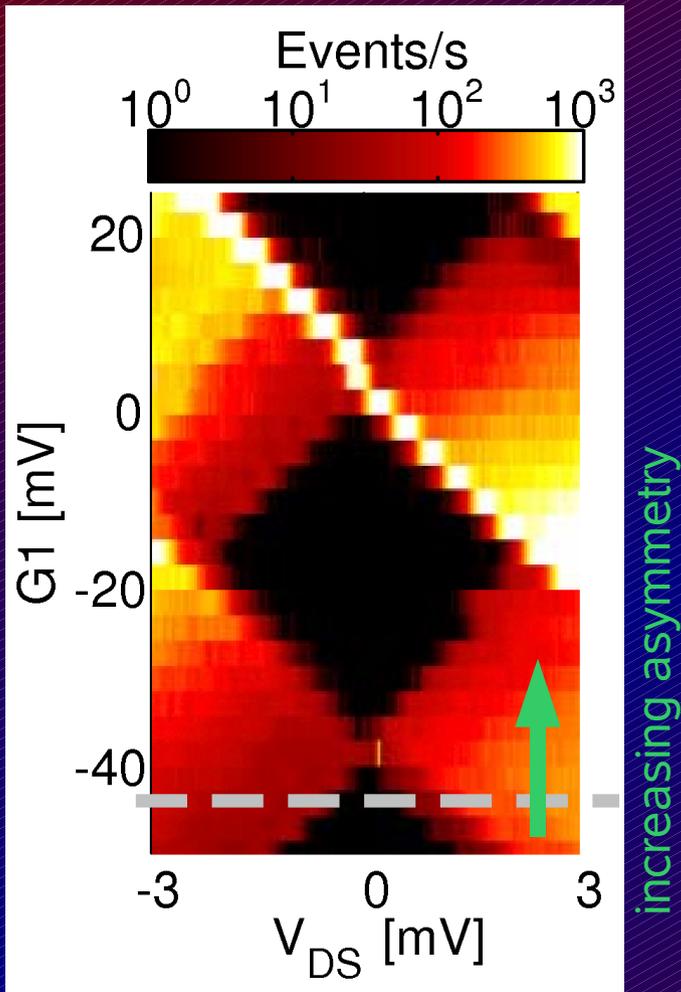
Theory: Hershfield *et al.*, PRB **47**, 1967 (1993)

Bagrets & Nazarov, PRB **67**, 085316 (2003)

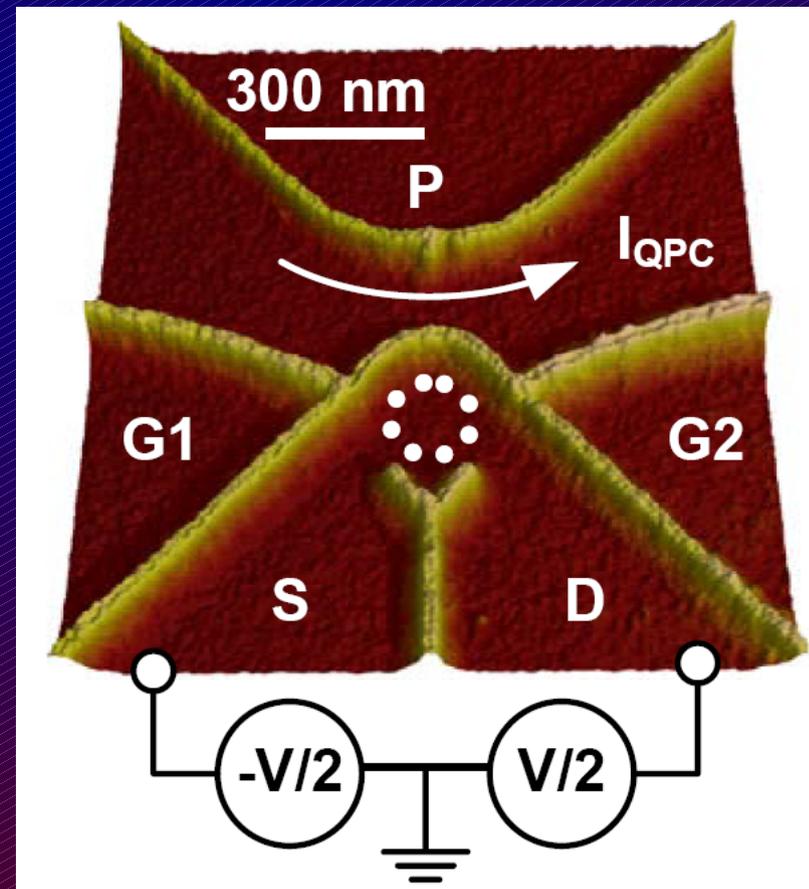
Bias dependence of the fluctuations



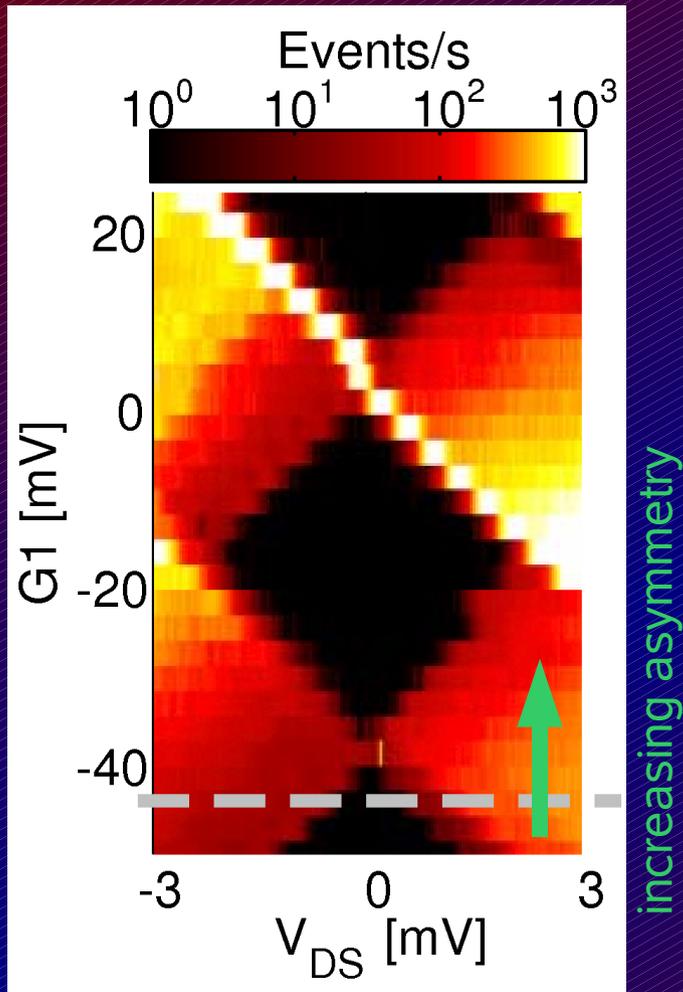
Adjustable asymmetry of the tunneling rates



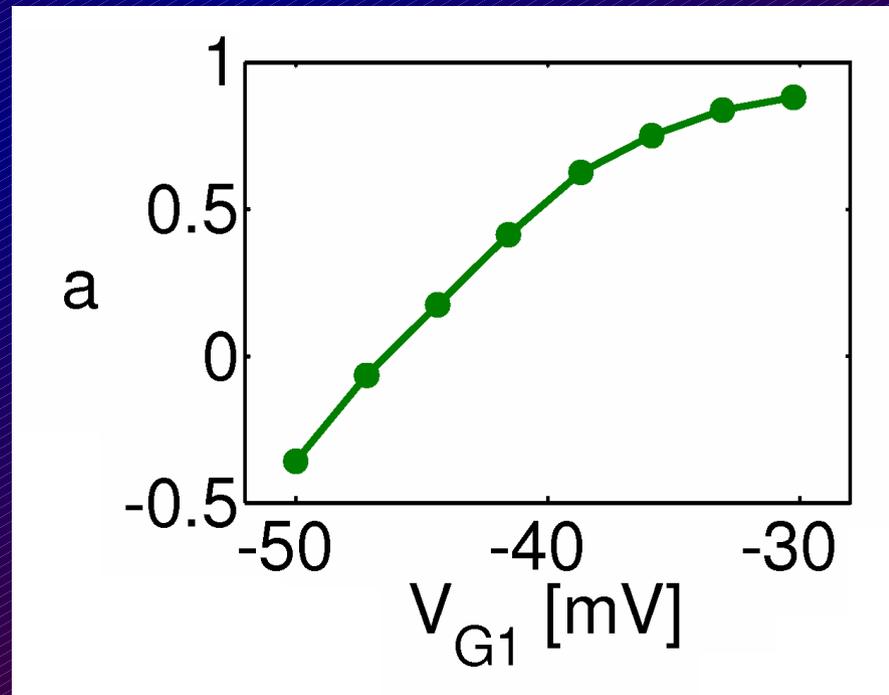
$$a = \frac{\Gamma_S - \Gamma_D}{\Gamma_S + \Gamma_D}$$



Adjustable asymmetry of the tunneling rates

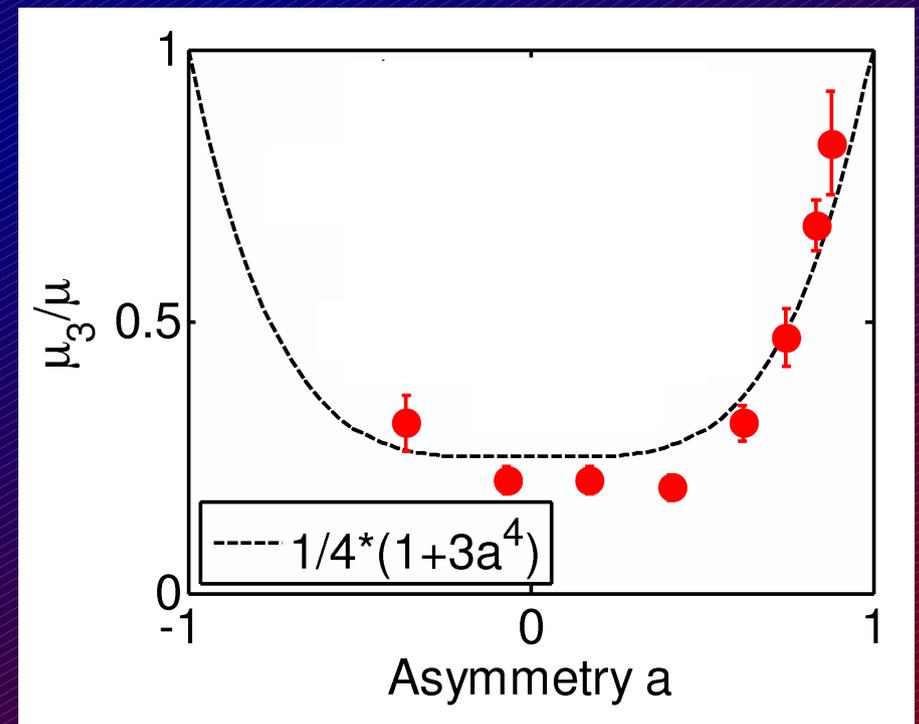
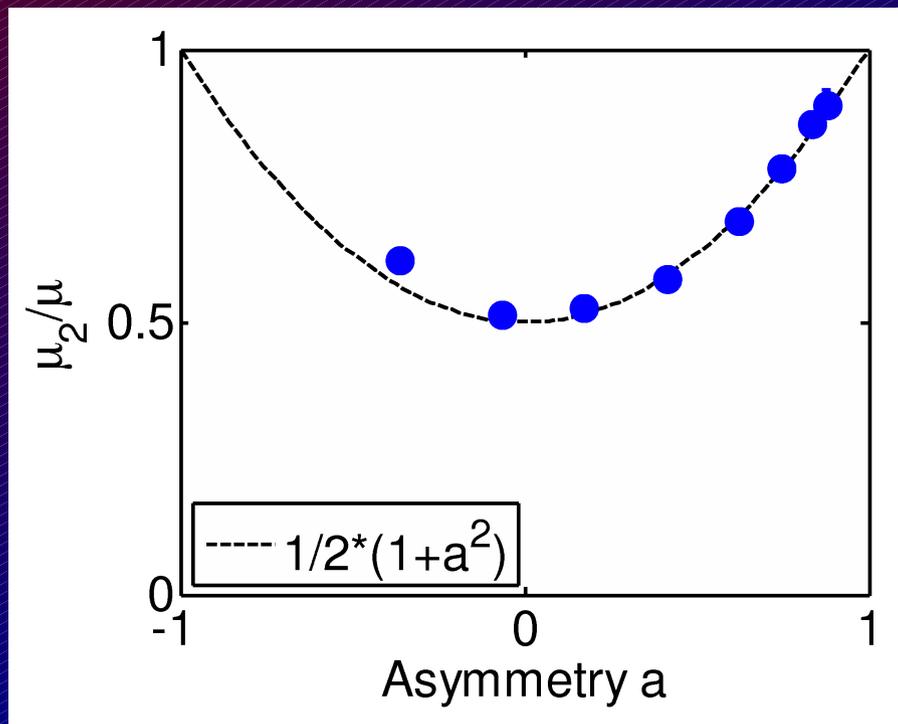


$$a = \frac{\Gamma_S - \Gamma_D}{\Gamma_S + \Gamma_D}$$



Current fluctuations vs. asymmetry

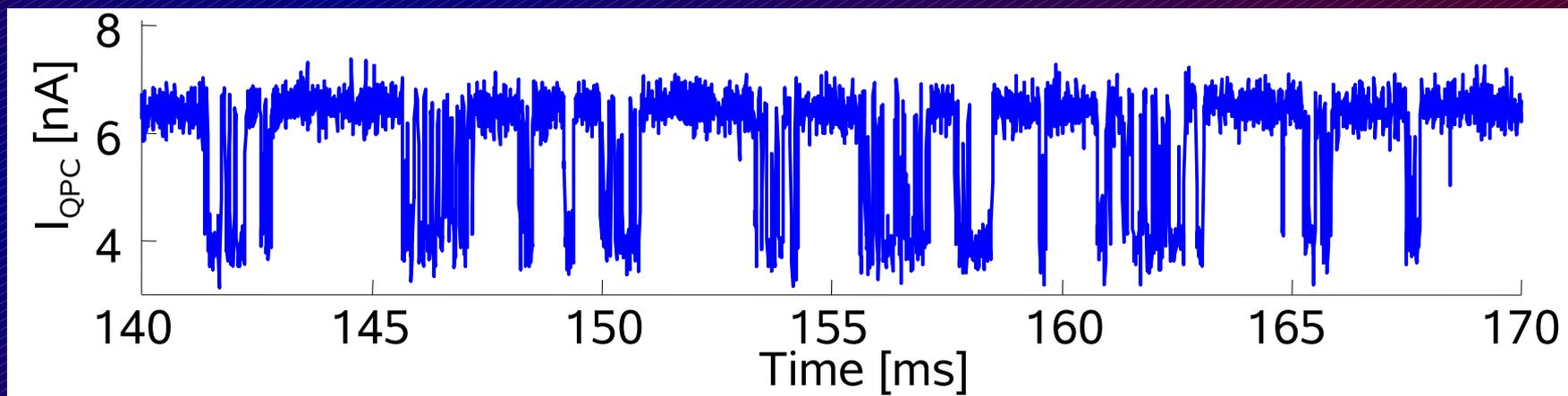
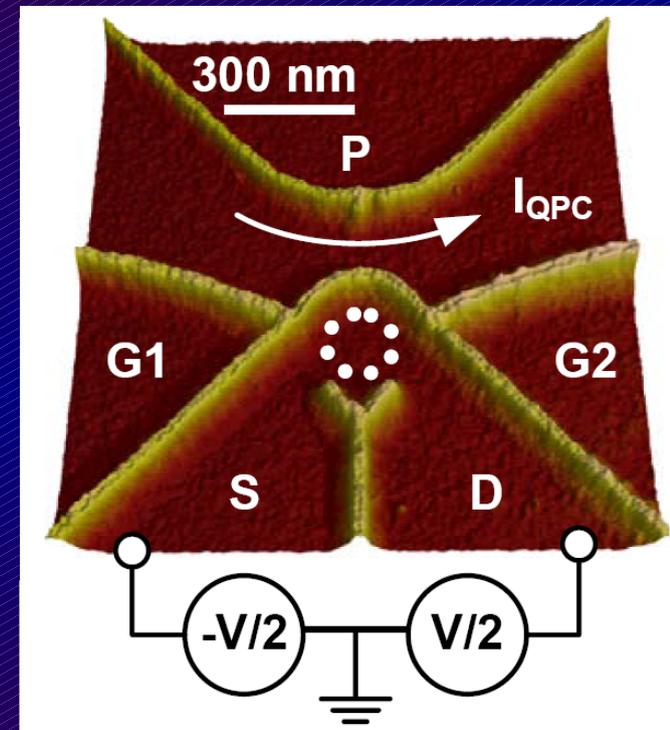
- Reduction of the second and third moments for symmetric coupling



Theory: Hershfield *et al.*, PRB 47, 1967 (1993)
Bagrets & Nazarov, PRB 67, 085316 (2003)

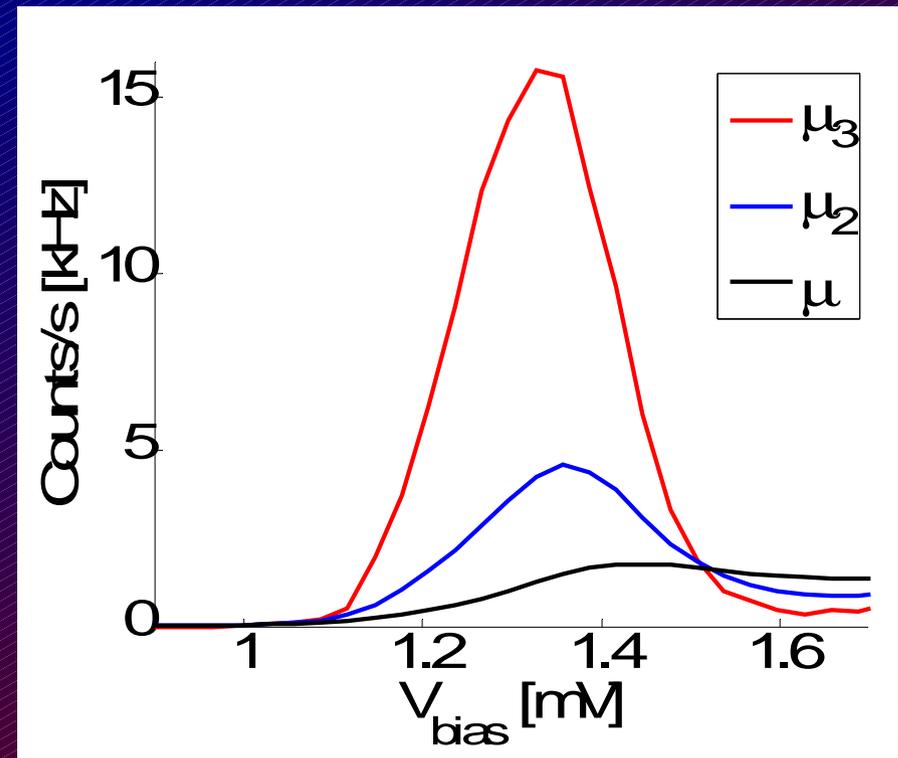
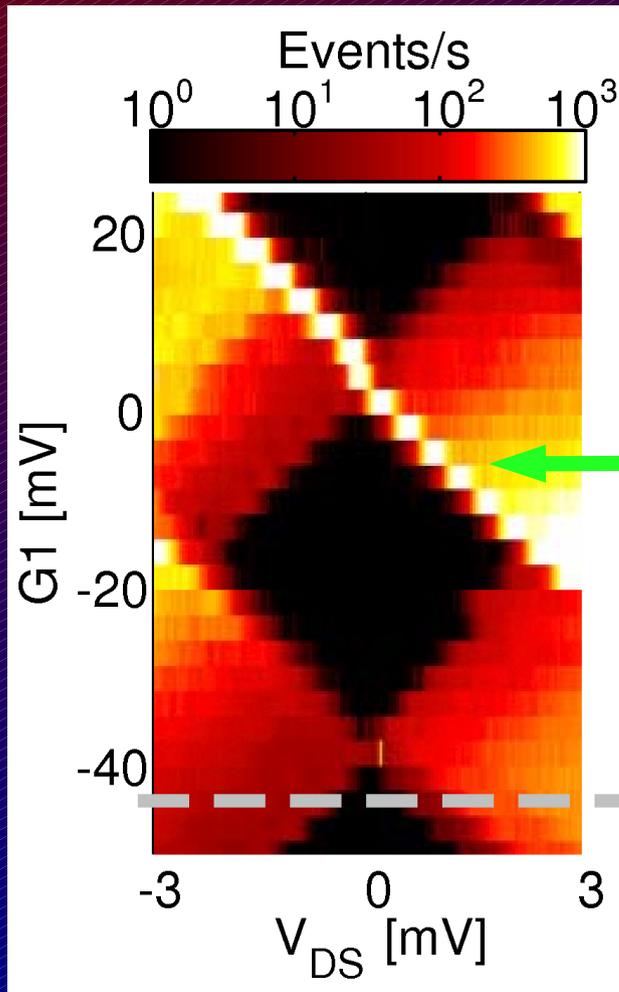
Bunching of electrons

- Two time scales
 - $\Gamma_1 \sim 20$ kHz, $\Gamma_0 \sim 1.5$ kHz
- Fast tunneling sometimes blocked by a slow tunneling



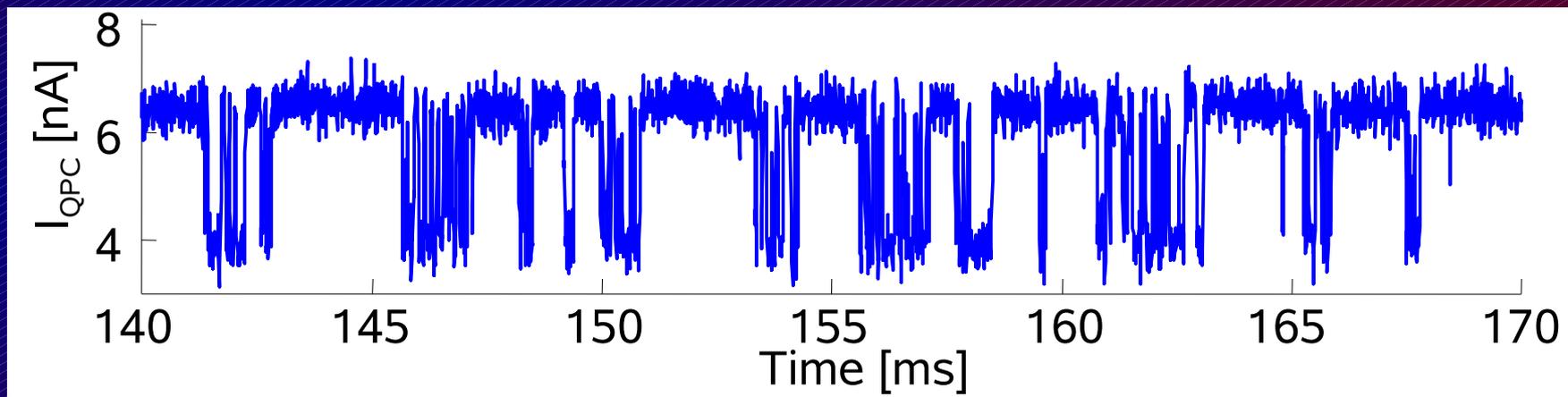
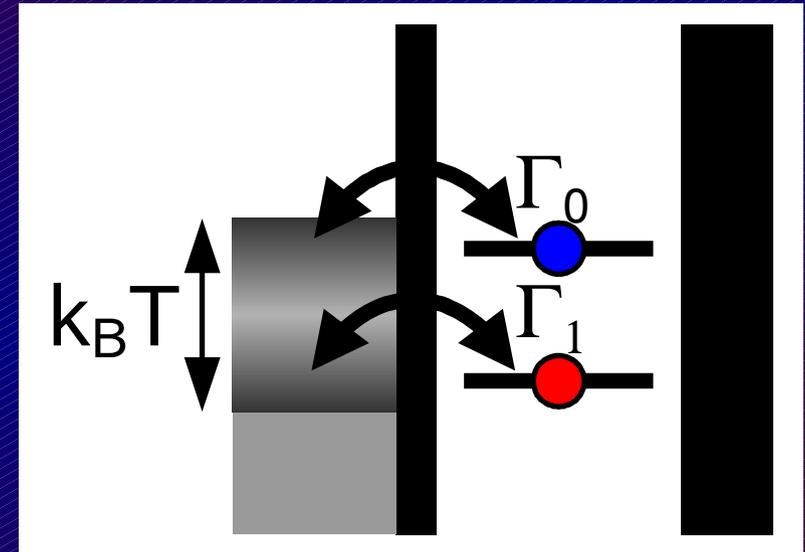
Bunching of electrons: super-poissonian shot noise

- super-poissonian noise occurs at the edge of conductance steps



Bunching of electrons: the model

- Needs two states with different coupling to the leads
 - the slow state blocks the conduction, due to Coulomb blockade
 - similar to Belzig, PRB 71, 161301(R) (2005)

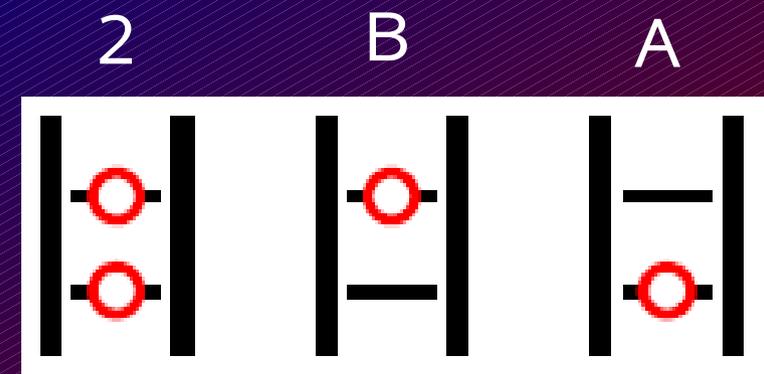
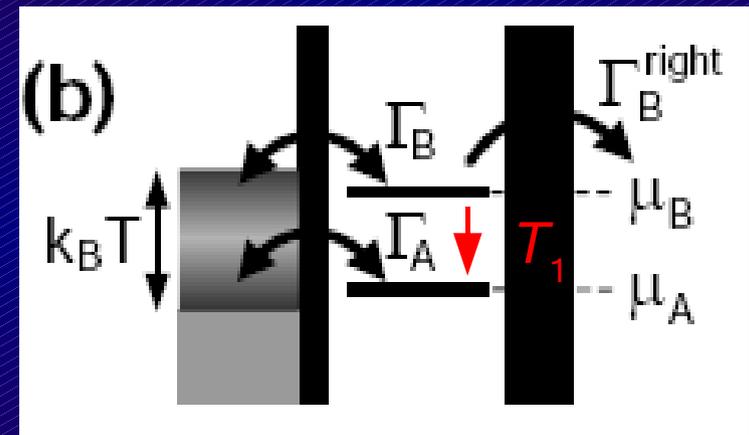


Bunching of electrons: the model

➤ Master equation

$$\frac{d}{dt} \begin{pmatrix} p_A \\ p_B \\ p_2 \end{pmatrix} = M \begin{pmatrix} p_A \\ p_B \\ p_2 \end{pmatrix}$$

$$M(\chi) = \begin{pmatrix} -\Gamma_B^i & \frac{1}{T_1} & (\Gamma_B^o + \Gamma_B^r) e^{i\chi} \\ 0 & -\Gamma_A^i - \frac{1}{T_1} & \Gamma_A^o e^{i\chi} \\ \Gamma_B^i & \Gamma_A^i & -\Gamma_A^o - \Gamma_B^o - \Gamma_B^r \end{pmatrix}$$



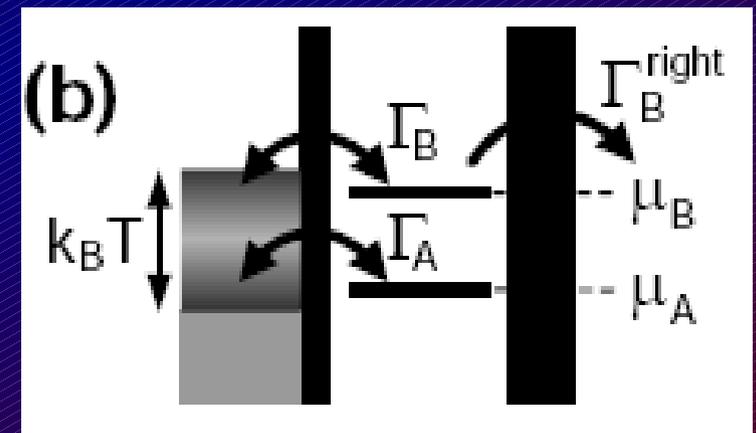
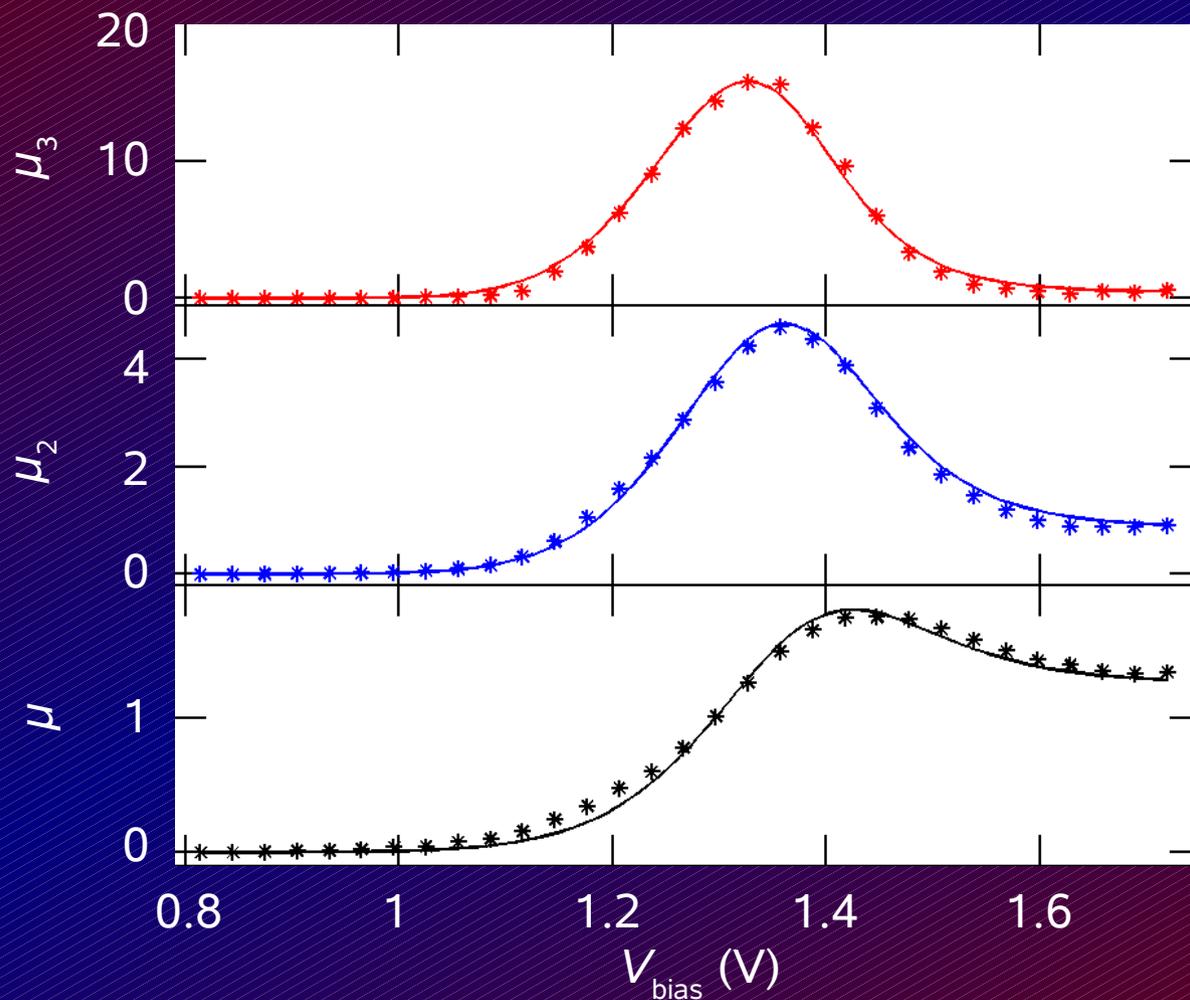
- Cumulant generating function determined from the eigenvalues of $M(\chi)$

Bagrets & Nazarov, PRB 67, 085316 (2003)

Belzig, PRB 71, 161301(R) (2005)

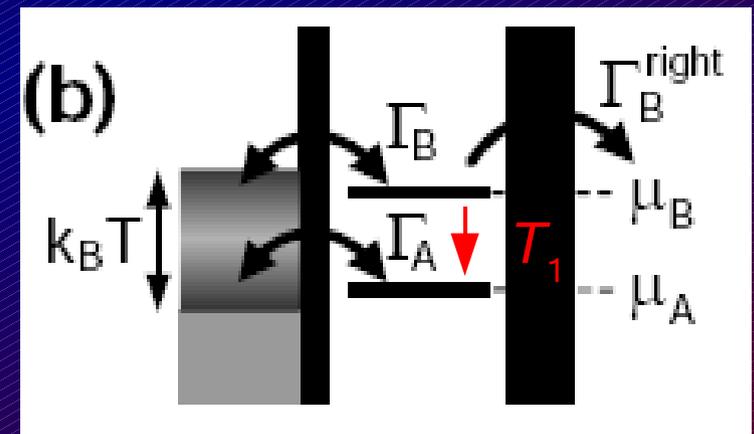
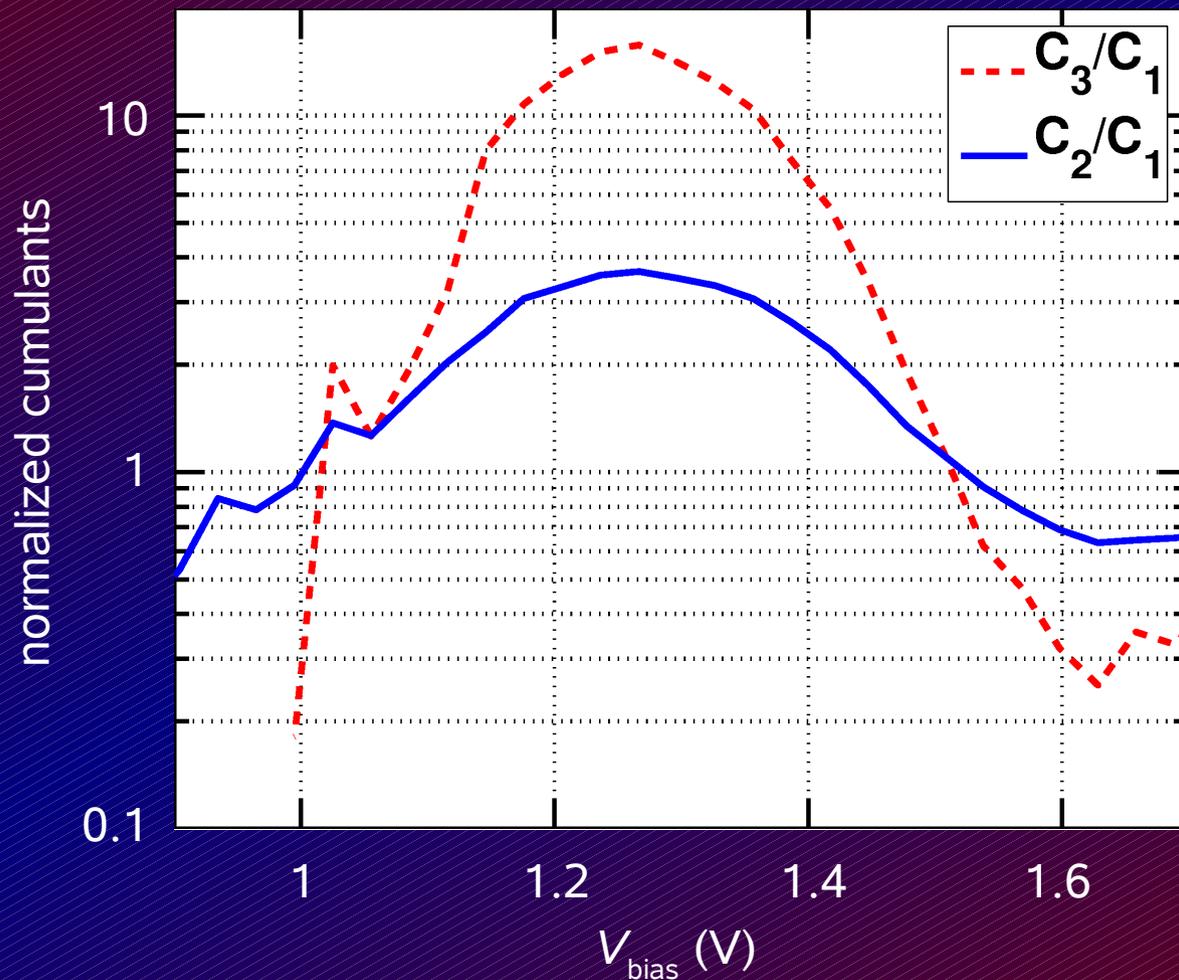
Bunching of electrons: super-poissonian shot noise

- Well described by our model



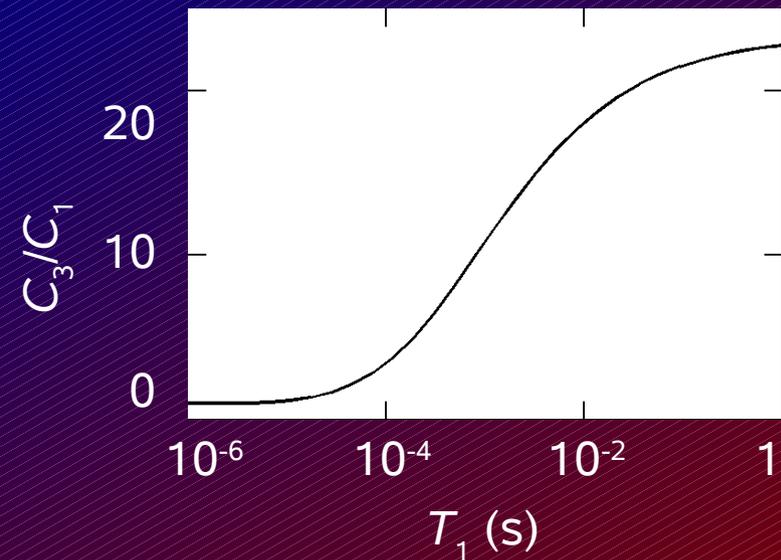
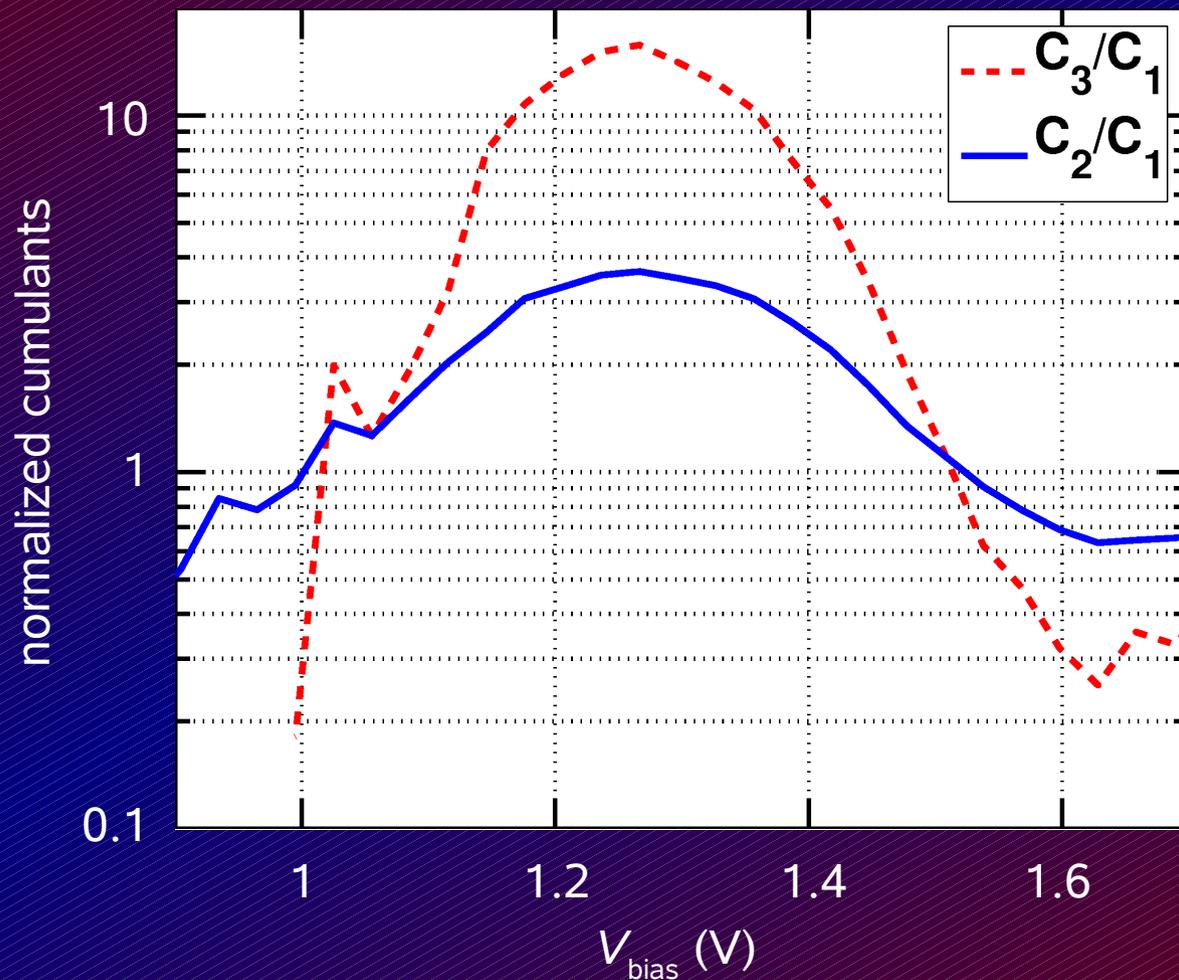
Bunching of electrons: super-poissonian shot noise

- Well described by our model, but requires a long relaxation rate: $T_1 > 1\text{ms}$... spin effect?



Bunching of electrons: super-poissonian shot noise

- Well described by our model, but requires a long relaxation rate: $T_1 > 1\text{ms}$... spin effect?



Conclusion



- Real-time detection of single electron traveling through a semiconductor quantum dot
- Measurement of current fluctuations
- Reduction of both the second and the third moments for symmetric coupling
- Bunching of electrons due to Coulomb blockade (information about relaxation time)
- **Noise measurements are now available in lateral quantum dots**
(even full counting statistics)